

FS 2024/25

MSE-422 – Advanced Metallurgy

11-Thermodynamic and kinetic modeling and simulation

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- The CALPHAD method is based on the thermodynamic modeling of solution phases and stoichiometric compounds
- For all solution phases, the Gibbs energy G must be determined

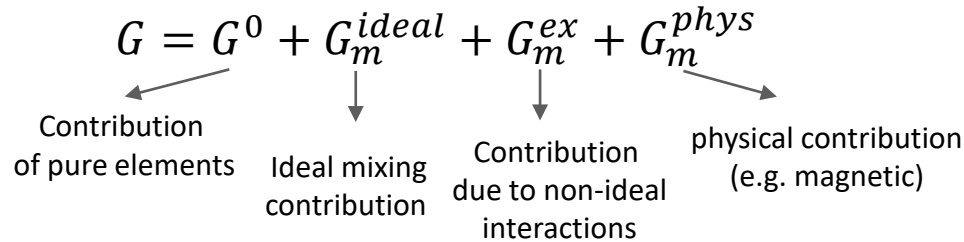
$$G = G^0 + G_m^{ideal} + G_m^{ex} + G_m^{phys}$$


Diagram illustrating the components of the Gibbs energy equation:

- G^0 : Contribution of pure elements
- G_m^{ideal} : Ideal mixing contribution
- G_m^{ex} : Contribution due to non-ideal interactions
- G_m^{phys} : physical contribution (e.g. magnetic)

- Solution phase = solid solution (e.g. FCC, BCC, HCP....) or LIQUID
- Thermodynamic modeling implies the mathematical description of each phase, thereby considering the physical properties of the phase (e.g. lattice structure, C_p , etc.) as precisely as possible

Thermodynamic modeling

Pure elements

- The integral Gibbs energy of a pure element φ in the state i (neglecting p) is

$$G_i^\varphi(T) = H - TS$$

- Considering the definitions of H and S this can be written as

$$G_i^\varphi(T) = H_m^{SER} + \int_0^T C_p dT - T \int_0^T \frac{C_p}{T} dT$$

with H_m^{SER} : enthalpy of the element/substance in its defined reference state at 298.15K and 1bar

- In the case of elements, the functions recommended by Scientific Group Thermodata Europe (SGTE) are generally used for representing G_i^φ :

$${}^0G_i^\varphi(T) = G_i^\varphi(T) - H_m^{SER} = a + bT + cT \ln(T) + dT^2 + eT^3 + fT^{-1} + gT^7 + hT^{-9}$$

Where a, b, c are the model parameters

- For ferromagnetic substances, a magnetical ordering term G_{mag}^φ needs to be added:

$$G_{mag}^\varphi(T) = RT \ln(1 + \beta) g(\tau)$$

where $\tau = T/T_C$, T_C is the critical temperature for magnetic transition, β is the magnetic moment in Bohr magneton and $g(\tau)$ is the magnetic ordering function

Stoichiometric compounds

- The same approach can be extended to model the Gibbs energy function of a stoichiometric compound θ

$$G_i^\theta(T) = \sum_i v_i H_i^{SER} + A + BT + CT \ln(T) + DT^2 + \dots$$

where A, B, C... are the model parameters and v_i are the stoichiometric coefficients for the elements that make up the compound

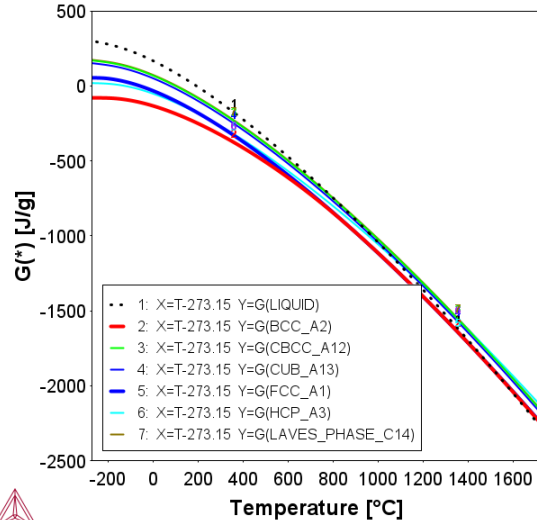
- For compounds with zero C_p of formation or where C_p is not known, a simpler expression can be used

$$G_i^\theta(T) = \sum_i v_i G_i^{ref} + \Delta_f H^\theta - T \Delta_f S^\theta = \sum_i v_i G_i^{ref} + I + JT$$

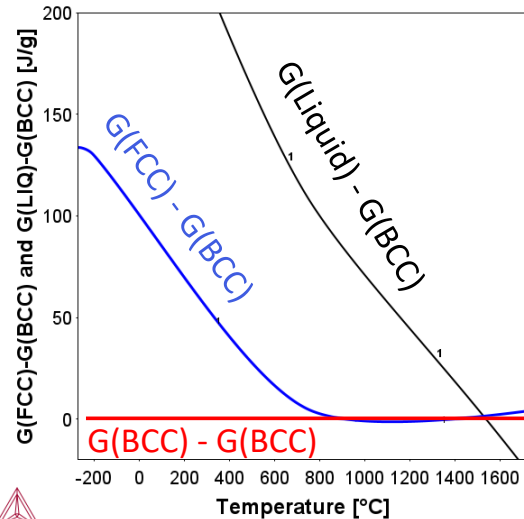
Thermodynamic modeling

Example: Gibbs energies of phases in pure iron

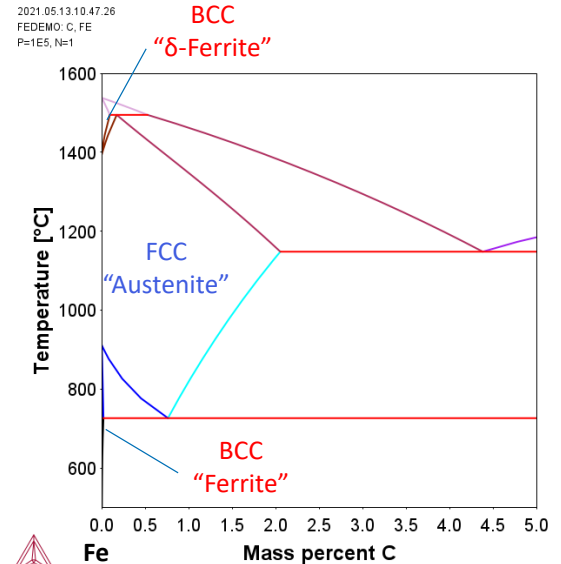
2021.05.13.10.40.56
FEDEMO: FE
P=1E5, B=1



2021.05.13.10.27.44
FEDEMO: FE
P=1E5, B=1



2021.05.13.10.47.26
FEDEMO: C, FE
P=1E5, N=1

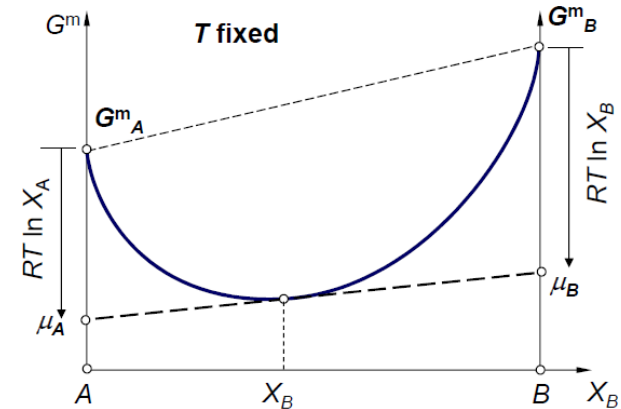
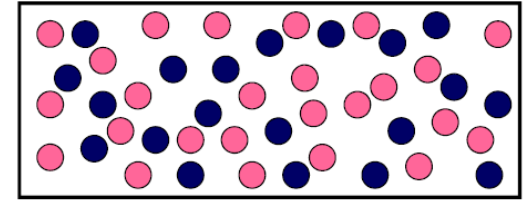


Thermodynamic modeling

Ideal solutions

- The simplest model for binary solutions is the ideal solution model
- It can readily be extended to solutions of higher order (more components)
- Its assumptions are that
 - the distribution of the atoms A and B is completely random
 - the exchange energy between atoms of type A and B is equal to the average exchange energy of A-A bonds and B-B bonds.

$$G_m^\alpha = \sum_i x_i G_i^0 + RT \sum_i x_i \ln x_i$$



Thermodynamic modeling

Regular solutions

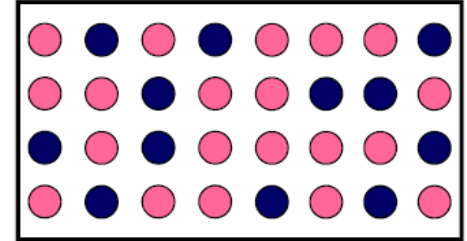
- In regular solutions the exchange energy in A-B bonds is no longer equal to the average of the bond energies of A-A and B-B bonds.
- A mixing enthalpy $\Delta H_{mix} = \Omega X_A X_B$ needs to be added
- Ω is a (temperature dependent) parameter that describes the interaction between A and B atoms

$$\Omega = N_a z \left(\varepsilon_{AB} - \left(\frac{\varepsilon_{AA} + \varepsilon_{BB}}{2} \right) \right)$$

N_a : Avogrado constant
 z : coordination number

$$G_m^\alpha = \sum_i x_i G_i^0 + RT \sum_i x_i \ln x_i + \sum_i \sum_{j>i} x_i x_j \Omega_{ij}$$

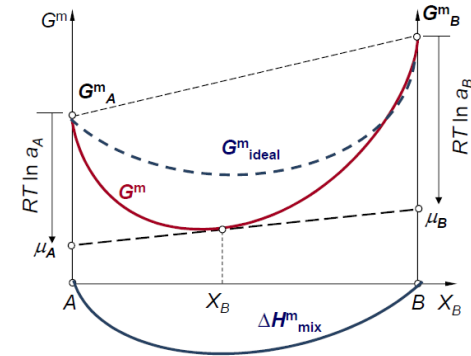
- Depending on the type of bonds, $\Omega < 0$ (A and B “like” each other) or $\Omega > 0$ (A and B “dislike” each other)



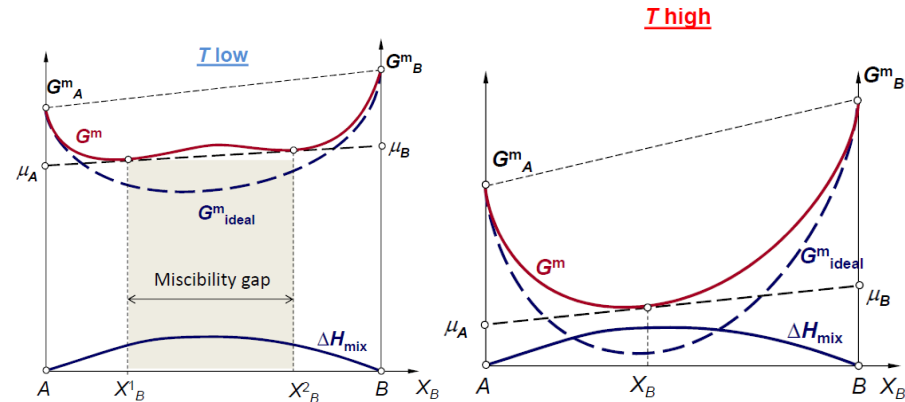
Thermodynamic modeling

Regular solutions

- Case $\Omega < 0$: A and B “like” each other
 - G_m remains concave at any temperature and the enthalpy of mixing only reinforces the “attractiveness” of A and B atoms.



- Case $\Omega > 0$: A and B “dislike” each other
 - At high temperature, $-T\Delta S_m$ still dominates and G_m remains concave. At lower temperature, the enthalpy of mixing dominates and can lead to demixing (miscibility gap).



Real solutions

- Real solutions deviate to variable extent from ideal or regular solutions. The reason for this can be that
 - 1) The exchange energy of A-B bonds is not independent of the composition (no unique Ω);
 - 2) The interaction between atoms of different type leads to a preferential arrangement of atoms A-B (instead of a random arrangement)
 - 3) Atoms can only take certain positions in a crystal or the liquid (e.g. for topological (size) or chemical reasons)
- These aspects are considered in
 - 1) The (sub)ⁿ-regular solution model (no physics, just numerical)
 - 2) The sub-lattice model
 - 3) The quasi-chemical approach

Real solutions – the (sub)ⁿ-regular solution model

- The enthalpy of mixing is not necessarily symmetric and higher order terms allow introducing some degree of skewedness to the enthalpy of mixing curve
- The Gibbs free energy is give as

$$G_m^\alpha = \sum_i x_i G_i^0 + RT \sum_i x_i \ln x_i + \sum_i \sum_{j>i} x_i x_j \sum_k L_{ij}^k (x_i - x_j)^k$$

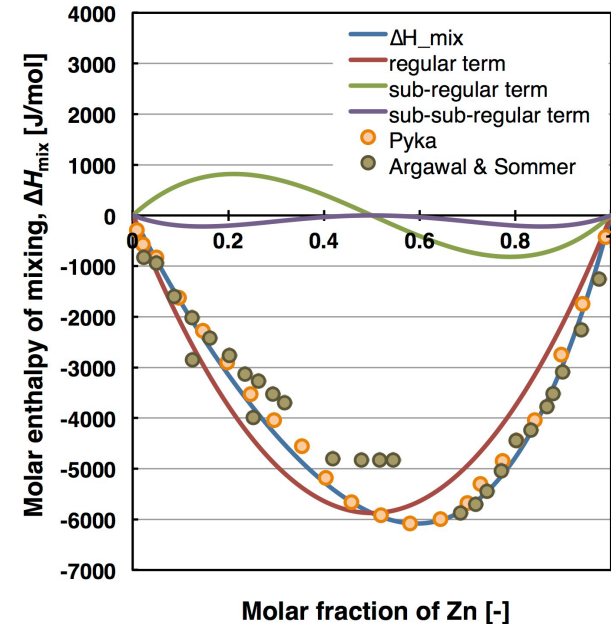
Where L_{ij}^k is a binary interaction parameter dependent on the value of k; $L_{ij}^k = A_k + B_k T + C_k T \ln T + \dots$

- This numerical approach, which facilitates the calculation of phase diagrams, is the Redlich-Kister formalism
- Examples for binary alloys of A and B
 - Regular solution: $L_{AB}^0 \neq 0, L_{AB}^k (k \geq 1) = 0$
 - Sub-regular solution: $L_{AB}^0, L_{AB}^1 \neq 0, L_{AB}^k (k \geq 2) = 0$
 - Sub-sub-regular solution: $L_{AB}^0, L_{AB}^1, L_{AB}^2 \neq 0, L_{AB}^k (k \geq 3) = 0$
- In some cases higher order interaction might not be neglected and a ternary interaction parameter is introduced: $G_{ABC}^{ex} = x_A x_B x_C L_{ABC}^k$

Thermodynamic modeling

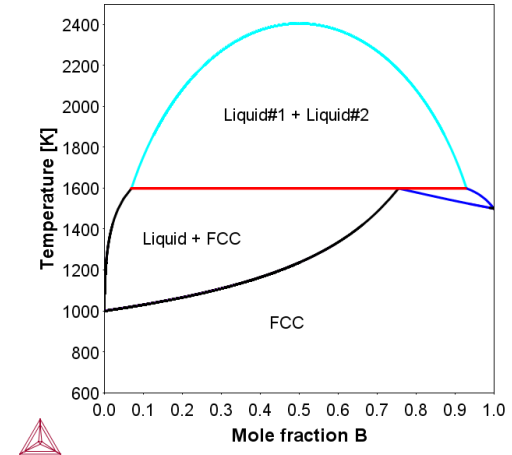
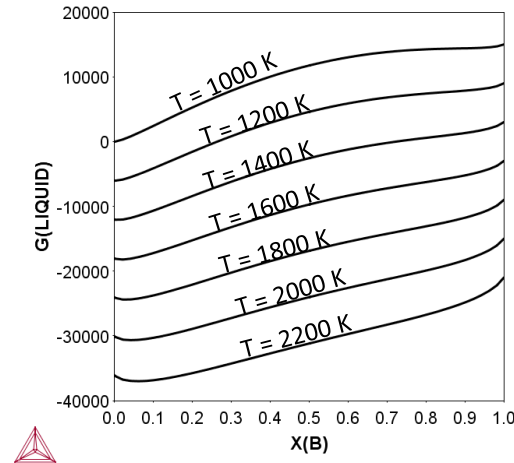
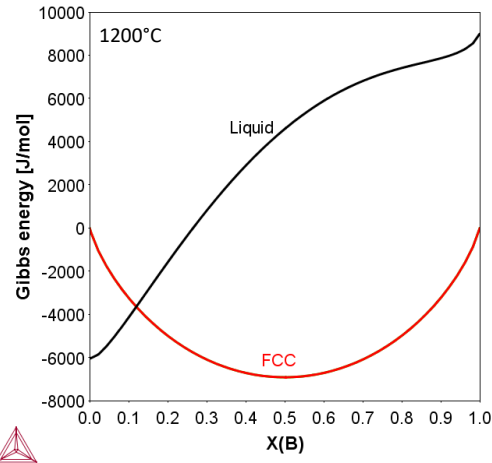
Real solutions – the (sub)ⁿ-regular solution model

- The sub-regular solution model is used for substitutional phases such as liquid, bcc, fcc, etc.
- It cannot be used for interstitial solutions, ordered intermetallics or ceramic compounds
- There is little evidence for the need of interaction parameters of any higher order than 2
- Prediction of thermodynamic properties of substitutional phases is based on binary and ternary terms
- Example: enthalpy of mixing of liquid Mg-Zn
 - experimental data from calorimetry measurements
 - Data fitting leads to
$$L_{Mg-Zn,liq}^0 = -23500 \text{ kJ/mol}$$
$$L_{Mg-Zn,liq}^1 = 8500 \text{ kJ/mol}$$
$$L_{Mg-Zn,liq}^2 = -3500 \text{ kJ/mol}$$



Thermodynamic modeling

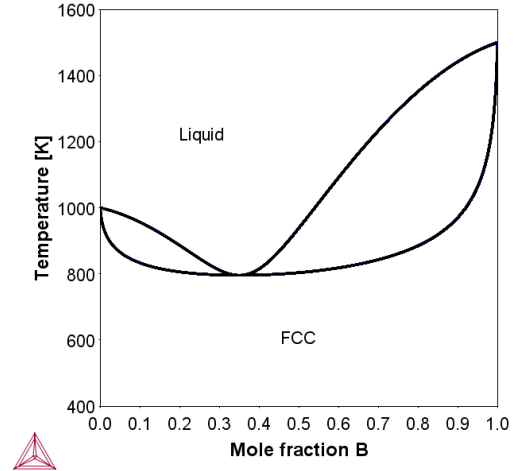
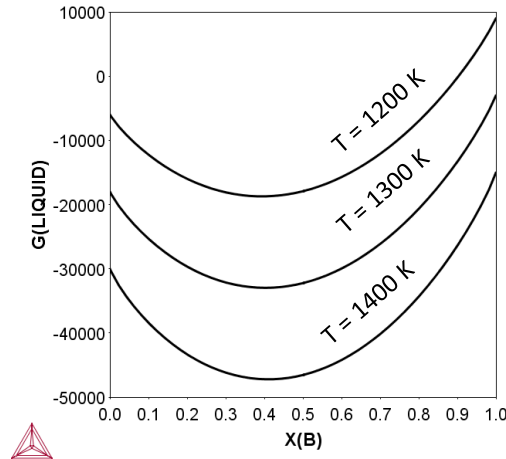
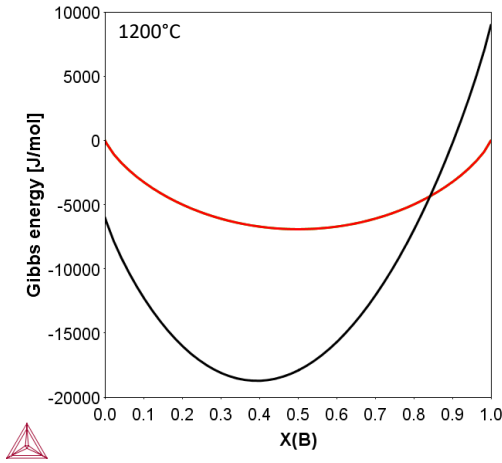
Example: influence of interaction parameters on phase stabilities



$$G_m^{ex,liquid} = x_A x_B {}^0L_{A,B} = x_A x_B \cdot (+40000) \rightarrow \text{negative excess term}$$

Thermodynamic modeling

Example: influence of interaction parameters on phase stabilities



$$G_m^{ex,liquid} = x_A x_B {}^0L_{A,B} = x_A x_B \cdot (-50000) \rightarrow \text{negative excess term}$$

Thermodynamic modeling

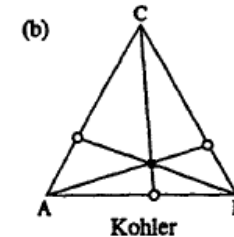
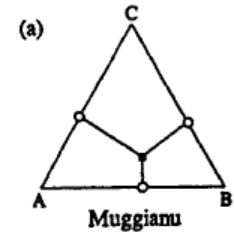
Real solutions – extrapolation to higher order systems

- Extrapolation of thermodynamic properties of alloys into multi-component systems is based on the summation of the binary and ternary excess parameters
- The formulae for doing this are based on various geometrical weightings of the mole fractions
- Muggianu's model:

$$G_{mix,ABC}^{ex} = x_A x_B \{L_{AB}^0 + L_{AB}^1 (x_A - x_B)\} \\ + x_B x_C \{L_{BC}^0 + L_{BC}^1 (x_B - x_C)\} \\ + x_A x_C \{L_{AC}^0 + L_{AC}^1 (x_A - x_C)\}$$

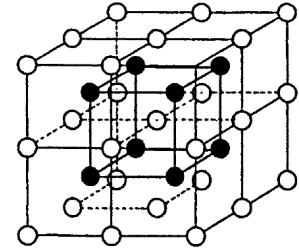
- Kohler's model:

$$G_{mix,ABC}^{ex} = (x_A - x_B)^2 \frac{x_A}{x_A + x_B} \cdot \frac{x_B}{x_A + x_B} \left\{ L_{AB}^0 + L_{AB}^1 \left(\frac{x_A - x_B}{x_A + x_B} \right) \right\} \\ + (x_B - x_C)^2 \frac{x_B}{x_B + x_C} \cdot \frac{x_C}{x_B + x_C} \left\{ L_{BC}^0 + L_{BC}^1 \left(\frac{x_B - x_C}{x_B + x_C} \right) \right\} \\ + (x_A - x_C)^2 \frac{x_A}{x_A + x_C} \cdot \frac{x_C}{x_A + x_C} \left\{ L_{AC}^0 + L_{AC}^1 \left(\frac{x_A - x_C}{x_A + x_C} \right) \right\}$$



Thermodynamic modeling

Real solutions – the sublattice model



- A sublattice (SL) phase is conceptualized as a combination of interlocking sublattices where various components can mix
- The model is phenomenological, focusing on macroscopic behavior, and does not inherently define any specific crystal structure (but can)
- Consider a phase with N distinct lattice sites, each forming a SL; the fractional site occupation of components on the different sublattices are:

$$y_i^s = \frac{n_i^s}{N^s} \text{ or including vacancies } y_i^s = \frac{n_i^s}{n_{Va}^s + \sum_i n_i^s}$$

Where n_i^s number of atoms of component i on sublattice s; N^s total number of sites on the sublattice

- Mole fractions are related to site fractions according to

$$x_i = \frac{\sum_s N^s y_i^s}{\sum_s N^s (1 - y_{Va}^s)}$$

- The ideal entropy of mixing is made up of the configurational contributions from mixing on each SL

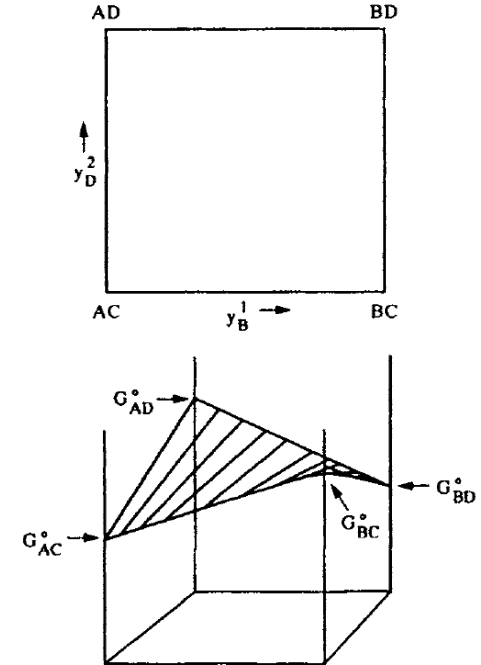
$$G_m^{id} = -TS_m^{id} = RT \sum_s N^s \sum_i y_i^s \ln y_i^s$$

Thermodynamic modeling

Real solutions – the sublattice model

- Lattice occupation:
 - Consider a sublattice phase consisting of the elements A, B, C, and D
 - A and B share the same sublattice and C and D are confined to another sublattice: $(A,B)_1(C,D)_1$
- End members
 - Four "complete occupation" configurations exist, representing pure combinations: pure A or B on SL1 and either pure C or D on SL2
 - These four configurations are referred to as "end members" and define the reference Gibbs energy surface.
- The Gibbs energy of the phase is a combination of the end members, calculated as:

$$G_m^{ref} = y_A y_C G_{AC}^0 + y_B y_C G_{BC}^0 + y_A y_D G_{AD}^0 + y_B y_D G_{BD}^0$$



Thermodynamic modeling

Real solutions – the sublattice model

- Consider again the two-SL system $(A,B)_1(C,D)_1$
- The interactions A-C, A-D, B-C, B-D are controlled by the Gibbs energies of the compounds AC, AD, BC and BD
- Mixing on the SLs controls A-B and C-D interactions and

$$G_{mix}^{ex} = y_A^1 y_B^1 L_{A,B:*}^0 + y_C^1 y_D^1 L_{*:C,D}^0$$

- A sub-regular model can be introduced by making the interactions compositionally dependent on the site occupation in the other SL

$$G_{mix}^{ex} = y_A^1 y_B^1 y_C^2 L_{A,B:C}^0 + y_A^1 y_B^1 y_D^2 L_{A,B:D}^0 + y_C^1 y_D^1 y_A^2 L_{A:C,D}^0 + y_C^1 y_D^1 y_B^2 L_{B:C,D}^0$$

- It is possible to add some site fraction dependence to these parameters

$$\text{e.g. } L_{A,B:C}^0 = y_A^1 y_B^1 y_C^2 \sum_v L_{A,B:C}^v (y_A^1 - y_B^1)^v$$

Thermodynamic modeling

Applications of the sublattice model

- The sublattice model is one of the most predominant methods used to describe solution and compound phases, e.g. for
 - Line compounds in ternary systems
 - e.g. Laves phase $\text{Fe}_2\text{TiSi} - (\text{Fe})_2(\text{Ti})_1(\text{Si})_1$
 - Interstitial phases,
 - e.g. carbides – $(\text{Fe}, \text{Cr}, \text{Ni}, \dots)_n(\text{C}, \text{Va})_m$
 - Complex intermetallic compounds with significant variation in stoichiometry,
 - e.g. $\text{Fe}_3\text{Al} - (\text{Fe})_3(\text{Al}, \text{Va})_1$
 - Order-disorder transformations
 - e.g. $\text{Cu}_3\text{Au} - \text{disordered phase } (\text{Cu}, \text{Au})_1 ; \text{ordered phase } (\text{Cu})_3(\text{Au})_1$

Thermodynamic modeling

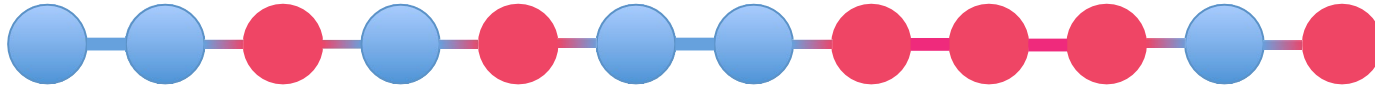
Real solutions – the quasi-chemical approach

- The regular solution model assumes a random distribution of atoms even though the enthalpy of mixing is not zero
- In reality, a random solution is only expected at very high temperatures when the entropy term overwhelms any tendency for ordering or clustering of atoms; the atom arrangement is no longer random and the entropy of mixing, $\Delta^{\text{mix}}S$, may differ from the ideal value
- The quasi-chemical approach takes into account that the interactions between atoms A and B are more or less energetic, which leads to preferential formation of A-B bonds ($\Omega < 0$) or A-A and B-B bonds ($\Omega > 0$)
- The model is so-called because it has a mass-action equation which is typical in chemical reaction theory

Thermodynamic modeling

Real solutions – the quasi-chemical approach

- Consider a binary system A-B of N atoms with the concentration x_A and x_B
- To understand the principle, the situation where the atoms are arranged on a 1-dimensional lattice is considered (“Ising”-model)



- The total number of bonds emanating from an atom of type A are $x_A N z$, with Z the coordination number of an atom ($Z=2$ in the 1D case)
- Of all the bonds emanating from an atom of type A, bonds of type A-B and of type B-A have to be subtracted (the A-B bonds are not interchangeable with the B-A bonds)
 - Number of A-A bonds: $N_{AA} = \frac{1}{2}(x_A Z N - N_{AB} - N_{BA})$
 - Number of B-B bonds: $N_{BB} = \frac{1}{2}(x_B Z N - N_{AB} - N_{BA})$

Thermodynamic modeling

The quasi-chemical approach – enthalpy term

- The total number of heterogeneous bonds, N_{het} , is $N_{\text{het}} = N_{AB} + N_{BA}$ and the number of A-B bonds is equal to the number of B-A bonds: $N_{AB} = N_{BA}$; The total number of bonds is $NZ/2$
- The concentration of bonds, n_{ik} , of the different types is:

$$\left. \begin{aligned} n_{AB} &= \frac{N_{AB}}{\frac{NZ}{2}} \\ n_{BA} &= \frac{N_{BA}}{\frac{NZ}{2}} \end{aligned} \right\} n_{\text{het}} = \frac{N_{AB} + N_{BA}}{\frac{NZ}{2}}$$

$$n_{AA} = \frac{N_{AA}}{\frac{NZ}{2}} = \left(x_A - \frac{1}{2} \frac{N_{AB}}{\frac{NZ}{2}} - \frac{1}{2} \frac{N_{BA}}{\frac{NZ}{2}} \right) = \left(x_A - \frac{1}{2} n_{AB} - \frac{1}{2} n_{BA} \right) = \left(x_A - \frac{n_{\text{het}}}{2} \right)$$

$$n_{BB} = \frac{N_{BB}}{\frac{NZ}{2}} = \left(x_B - \frac{1}{2} \frac{N_{AB}}{\frac{NZ}{2}} - \frac{1}{2} \frac{N_{BA}}{\frac{NZ}{2}} \right) = \left(x_B - \frac{1}{2} n_{AB} - \frac{1}{2} n_{BA} \right) = \left(x_B - \frac{n_{\text{het}}}{2} \right)$$

- These bonds contribute to the Gibbs free energy with γ_{AA} , γ_{BB} , and γ_{AB} (energy from the bonds) for A-A, B-B, and A-B as well as B-A bonds, respectively. The total Gibbs free energy, G , is then

$$G(N) = \frac{NZ}{2} (n_{AA}\gamma_{AA} + n_{\text{het}}\gamma_{AB} + n_{BB}\gamma_{BB}) - TS_c = \frac{NZ}{2} \left(\left(x_A - \frac{n_{\text{het}}}{2} \right) \gamma_{AA} + n_{\text{het}}\gamma_{AB} + \left(x_B - \frac{n_{\text{het}}}{2} \right) \gamma_{BB} \right) - TS_c$$

$$G(N) = \underbrace{\frac{NZ}{2} (x_A\gamma_{AA} + x_B\gamma_{BB})}_{x_A G_A(N) + x_B G_B(N)} + \underbrace{\frac{NZ}{2} \left(n_{\text{het}} \left(\gamma_{AB} - \left(\frac{\gamma_{AA} + \gamma_{BB}}{2} \right) \right) \right)}_{\Delta^{\text{mix}} G(N)} - TS_c$$

Thermodynamic modeling

The quasi-chemical approach – entropy term

- Assume that there is no contribution to entropy by e.g. change in lattice vibrations due to the A-B bonds $\left(\gamma_{AB} - \frac{\gamma_{AA} + \gamma_{BB}}{2}\right) = \left(\varepsilon_{AB} - \frac{\varepsilon_{AA} + \varepsilon_{BB}}{2}\right) = \omega$

- The total configurational entropy of the system of bonds, $S_{c,bonds}$, can be calculated as

$$S_{c,bonds} = k_B \ln W; \quad W = \frac{\frac{NZ}{2}!}{N_{AA}! N_{BB}! \frac{1}{2} N_{AB}! \frac{1}{2} N_{AB}!} = \frac{\frac{NZ}{2}!}{\left(n_{AA} \frac{NZ}{2}\right)! \left(n_{BB} \frac{NZ}{2}\right)! \left(n_{AB} \frac{NZ}{2}\right)! \left(n_{BA} \frac{NZ}{2}\right)!}$$

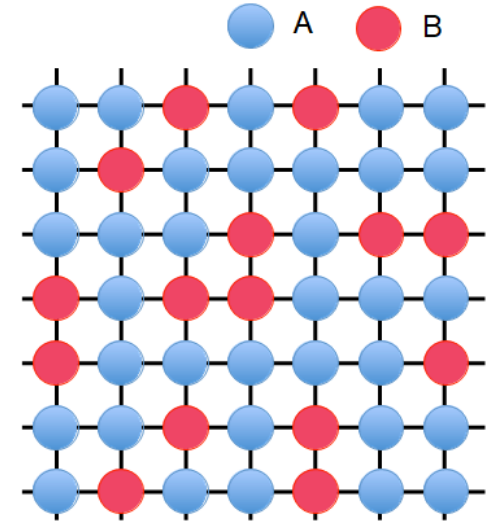
- With Sterling's formula, i.e. $\ln(x!) = x \ln(x)$, this can be written as:

$$\begin{aligned} S_{c,bonds} &= k_B \left[\left(\frac{NZ}{2}\right) \ln\left(\frac{NZ}{2}\right) - \left(n_{AA} \frac{NZ}{2}\right) \ln\left(n_{AA} \frac{NZ}{2}\right) - \left(n_{BB} \frac{NZ}{2}\right) \ln\left(n_{BB} \frac{NZ}{2}\right) - \left(n_{AB} \frac{NZ}{2}\right) \ln\left(n_{AB} \frac{NZ}{2}\right) - \left(n_{BA} \frac{NZ}{2}\right) \ln\left(n_{BA} \frac{NZ}{2}\right) \right] \\ S_{c,bonds} &= \frac{NZ}{2} k_B \left[\ln\left(\frac{NZ}{2}\right) - n_{AA} \ln\left(n_{AA} \frac{NZ}{2}\right) - n_{BB} \ln\left(n_{BB} \frac{NZ}{2}\right) - n_{AB} \ln\left(n_{AB} \frac{NZ}{2}\right) - n_{BA} \ln\left(n_{BA} \frac{NZ}{2}\right) \right] \\ S_{c,bonds} &= \frac{NZ}{2} k_B \left[\ln\left(\frac{NZ}{2}\right) - n_{AA} \left(\ln(n_{AA}) + \ln\left(\frac{NZ}{2}\right)\right) - n_{BB} \left(\ln(n_{BB}) + \ln\left(\frac{NZ}{2}\right)\right) - n_{AB} \left(\ln(n_{AB}) + \ln\left(\frac{NZ}{2}\right)\right) - n_{BA} \left(\ln(n_{BA}) + \ln\left(\frac{NZ}{2}\right)\right) \right] \\ S_{c,bonds} &= \frac{NZ}{2} k_B \left[\ln\left(\frac{NZ}{2}\right) (1 - n_{AA} - n_{BB} - n_{AB} - n_{BA}) - n_{AA} \ln(n_{AA}) - n_{BB} \ln(n_{BB}) - n_{AB} \ln(n_{AB}) - n_{BA} \ln(n_{BA}) \right] \\ S_{c,bonds} &= -\frac{NZ}{2} k_B \left[n_{AA} \ln(n_{AA}) + n_{BB} \ln(n_{BB}) + n_{AB} \ln(n_{AB}) + n_{BA} \ln(n_{BA}) \right] \end{aligned}$$

Thermodynamic modeling

The quasi-chemical approach

- In calculating the configurational entropy of the bond system, it was considered that A-B bonds, A-A bonds, B-B bonds, and B-A bonds are permutable amongst themselves
- However, A-B bonds are different from B-A bonds (not in their energy but with regard to permutability)
- It was assumed that the bonds can be placed randomly on the lattice of bonds, but there are already some restrictions on the Ising chain model
 - e.g. an A-A bond cannot be next to a B-B bond, neither can an A-B bond be to the left of a A-A bond (but a A-A bond or a B-A bond can).
- The arrangement of the bonds is hence somehow restricted (→ 2D model)
- Accounting for this in a 3D network of bonds in a detailed way becomes awfully complicated



Thermodynamic modeling

The quasi-chemical approach

- The entropy cannot be larger than in a fully random arrangement of the atoms.
- For 1 mole (i.e. $N = N_A$) of a binary mixture of atoms, the maximum configurational entropy is attained when $x_A = x_B = 0.5$ and $S_{c,atoms} = RZ\ln(2)$
- If we do the same calculation for the bonds, again for 1 mole of atoms, the maximum is obtained when $n_{AA} = n_{BB} = n_{AB} = n_{BA} = 0.25$, and its value is $S_{c,bonds} = RZ\ln(4)/2$
- It cannot be that the configurational entropy changes just because we looked at another entity, i.e. bonds rather than atoms
- the entropy calculated for the **bonds can be normalized** with regard the one obtained for the **atoms** by looking at their respective maximum value and assuming that also outside the maximum with the normalization factor Φ

$$\Phi = \frac{S_{c,atoms}}{S_{c,bonds}} = \frac{2\ln(2)}{Z\ln(4)} = \frac{1}{Z}$$

- The molar Gibbs free energy of mixing G_{mix} , as a function of the bond concentration of heterogeneous bonds

Thermodynamic modeling

The quasi-chemical approach

- The molar Gibbs free energy of mixing G_{mix} , as a function of the bond concentration of heterogenous bonds becomes

$$\Delta^{\text{mix}} g = \frac{N_A Z}{2} n_{\text{het}} \omega + T \frac{N_A Z}{2} \Phi k_B \left[\underbrace{\left(x_A - \frac{n_{\text{het}}}{2}\right)}_{n_{AA}} \ln \underbrace{\left(x_A - \frac{n_{\text{het}}}{2}\right)}_{n_{AA}} + \underbrace{\left(x_B - \frac{n_{\text{het}}}{2}\right)}_{n_{BB}} \ln \underbrace{\left(x_B - \frac{n_{\text{het}}}{2}\right)}_{n_{BB}} + \underbrace{\frac{n_{\text{het}}}{2}}_{n_{AB}} \ln \underbrace{\left(\frac{n_{\text{het}}}{2}\right)}_{n_{AB}} + \underbrace{\frac{n_{\text{het}}}{2}}_{n_{BA}} \ln \underbrace{\left(\frac{n_{\text{het}}}{2}\right)}_{n_{BA}} \right]$$

- The equilibrium value for the bond concentration n_{het} is found when G_{mix} is minimum
- Setting the 1st derivative 0, rearranging and considering $\Phi=1/Z$

$$TR \left[\ln \left(x_A - \frac{n_{\text{het}}}{2} \right) + \ln \left(x_B - \frac{n_{\text{het}}}{2} \right) - 2 \ln \left(\frac{n_{\text{het}}}{2} \right) \right] = 2N_A Z \omega$$

$$\text{with } N_A Z \omega = \Omega: \quad \exp \left(\frac{2\Omega}{RT} \right) = \frac{\left(x_A - \frac{n_{\text{het}}}{2} \right) \left(x_B - \frac{n_{\text{het}}}{2} \right)}{2 \left(\frac{n_{\text{het}}}{2} \right)} = \frac{[n_{AA}][n_{BB}]}{[n_{AB}][n_{BA}]}$$

Looks like mass-action:



Thermodynamic modeling

Applications of the quasi-chemical approach

- The quasi-chemical approach is particularly effective in modeling non-ideal systems where interactions between components lead to significant short-range ordering, clustering, or other non-random atomic arrangements.
 - Description of short-range order in Cu-Zn or Ni-Cr alloys
 - Description of miscibility gaps in the Ni-Cr and Fe-Cr systems
 - Description of ionic systems, such as molten salts or ceramics

Thermodynamic modeling

Example: assessment of the binary Ge-Ni system

CALPHAD: Computer Coupling of Phase Diagrams and Thermochemistry 38 (2012) 23–34



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Thermodynamic study and re-assessment of the Ge-Ni system

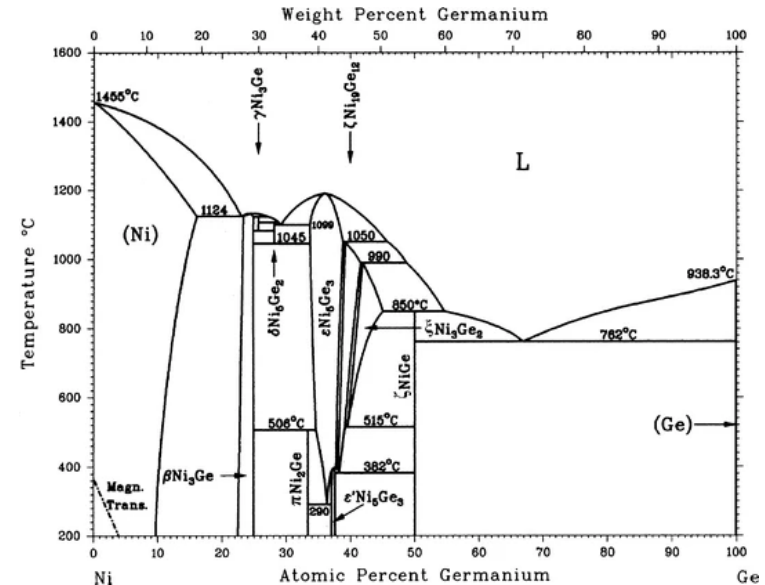
Shan Jin^{a,b,*}, Christian Leinenbach^a, Jiang Wang^a, Liliana I. Duarte^a, Simona Delsante^c, Gabriella Borzone^c, Andrew Scott^d, Andrew Watson^d

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Thermodynamic modeling

Example: assessment of the binary Ge-Ni system

■ Work flow for thermodynamic assessment

- Review of literature data
- Own measurements (phase diagram, calorimetry)
- Thermodynamic modeling

- Modeling of pure elements ${}^0G_i^\varphi(T) = G_i^\varphi(T) - H_i^{SER} = a + bT + cT \ln T + dT^2 + eT^3 + fT^{-1} + gT^7 + hT^{-9}$

- Modeling of solution phases $G_m^\varphi = \sum x_i {}^0G_i^\varphi + RT \sum x_i \ln(x_i) + {}^E G_m^\varphi + {}^{mag} G_m^\varphi$

$${}^E G_m^\varphi = X_{Ge} X_{Ni} \sum_{j=0}^n {}^{(j)} L_{Ge,Ni}^\varphi (X_{Ge} - X_{Ni})^j \quad {}^{(j)} L_{Ge,Ni}^\varphi = A_j + B_j T + C_j T \ln T \quad {}^{mag} G_m^\varphi = RT \ln(\beta^\varphi + 1) f(\tau^\varphi)$$

- Modeling of ordered phases $G_m = G_m^{disord}(x_i) + G_m^{ord}(y'_i, y''_i) - G_m^{ord}(x_i)$

- Modeling of B8-type non-stoichiometric compound Ni_5Ge_3
→ 3-sublattice model (Ge)(Ni)(Va,Ni)

$$G_m^{Ni_5Ge_3} = y_{Va}^{III} {}^0G_{Ge:Ni:Va}^{Ni_5Ge_3} + y_{Ni}^{III} {}^0G_{Ge:Ni:Ni}^{Ni_5Ge_3} + RT(y_{Ni}^{III} \ln y_{Ni}^{III} + y_{Va}^{III} \ln y_{Va}^{III}) + {}^E G_m^{Ni_5Ge_3}$$

$${}^0G_{Ge:Ni:Va}^{Ni_5Ge_3} = {}^0G_{Ge}^{DIAMOND_A4} + {}^0G_{Ni}^{FCC_A1} + E + FT$$

$${}^0G_{Ge:Ni:Ni}^{Ni_5Ge_3} = {}^0G_{Ge}^{DIAMOND_A4} + 2 {}^0G_{Ni}^{FCC_A1} + E' + F'T$$

$${}^E G_m^{Ni_5Ge_3} = y_{Va}^{III} y_{Ni}^{III} {}^E G_{Ge:Ni:Va,Ni}^{Ni_5Ge_3} \quad {}^E G_{Ge:Ni:Va,Ni}^{Ni_5Ge_3} = \sum_{n=0}^n (G_n + H_n T) (y_{Ni}^{III} - y_{Va}^{III})^n$$

- Modeling of stoichiometric phases

$$G_m^{NiGe} = 0.5 {}^0G_{Ni}^{FCC_A1} + 0.5 {}^0G_{Ge}^{DIAMOND_A4} + a + b + cT \ln T + dT^2 + eT^3 + fT^{-1} + gT^7$$

$$G_m^{Ni_xGe_y} = \frac{x}{x+y} {}^0G_{Ni}^{FCC_A1} + \frac{y}{x+y} {}^0G_{Ge}^{DIAMOND_A4} + I_{Ni_xGe_y} + J_{Ni_xGe_y} T$$

Thermodynamic modeling

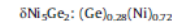
Example: assessment of the binary Ge-Ni system

■ Work flow for thermodynamic assessment

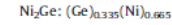
■ Parameter optimization (→ least-square fitting algorithm in Thermo-Calc)

Table 2
Thermodynamic parameters of the Ge-Ni binary system.^a

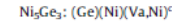
Phase	Thermodynamic parameters
Liquid: (Ge,Ni) ₁	${}^0L_{GeNi} = -167121.320 + 1557 - 157 \ln T$ ${}^1L_{GeNi} = 84737.489 - 25.014T$ ${}^2L_{GeNi} = 37441.590 - 16.001T$ ${}^3L_{GeNi} = -63650.323 + 21.983T$
FCC_A1 ^b : (Ge,Ni) ₁	${}^0F_{GeNi}^{FCC,A1} = -122000 + 36.88T$ ${}^1F_{GeNi}^{FCC,A1} = 134000 - 46.8T$ ${}^2F_{GeNi}^{FCC,A1} = -3750$
βNi_3Ge^b : (Ge,Ni) _{0.75} (Ge,Ni) _{0.25}	${}^0G_{GeNi}^{\beta Ni_3Ge} - 0.75 {}^0G_{Ge}^{DIAMOND,A4} - 0.25 {}^0G_{Ni}^{FCC,A1} = -46827.192 + 3.05T$ ${}^0G_{Ni_3Ge}^{\beta Ni_3Ge} - 0.25 {}^0G_{Ge}^{DIAMOND,A4} - 0.75 {}^0G_{Ni}^{FCC,A1} = -46827.192 + 3.05T$ ${}^0G_{GeGe}^{\beta Ni_3Ge} - 0.5 {}^0G_{Ge}^{DIAMOND,A4} = 0$ ${}^0G_{NiNi}^{\beta Ni_3Ge} - 0.5 {}^0G_{Ni}^{FCC,A1} = 0$ ${}^1G_{GeNi}^{\beta Ni_3Ge} = -93654.384 + 6.1T$ ${}^1G_{Ni_3Ge}^{\beta Ni_3Ge} = -93654.384 + 6.1T$ ${}^1G_{GeNi}^{\beta Ni_3Ge} = 23700 - 9.72T$ ${}^1G_{Ni_3Ge}^{\beta Ni_3Ge} = 23700 - 9.72T$ ${}^0G_{Ni_3Ge}^{\beta Ni_3Ge} = 0$ ${}^0G_{NiGe}^{\beta Ni_3Ge} = 0$ ${}^1G_{Ni_3Ge}^{\beta Ni_3Ge} = 7900 - 3.24T$ ${}^1G_{GeNi}^{\beta Ni_3Ge} = 7900 - 3.24T$
γNi_3Ge : (Ge) _{0.25} (Ni) _{0.75}	${}^0G_{Ni_3Ge}^{\gamma Ni_3Ge} - 0.25 {}^0G_{Ge}^{DIAMOND,A4} - 0.75 {}^0G_{Ni}^{FCC,A1} = -34315 + 4.301T$



$${}^0G_{Ni_5Ge_2}^{\delta Ni_5Ge_2} - 0.28 {}^0G_{Ge}^{DIAMOND,A4} - 0.72 {}^0G_{Ni}^{FCC,A1} = -34918 + 3.69T$$



$${}^0G_{Ni_2Ge}^{\delta Ni_2Ge} - 0.333 {}^0G_{Ge}^{DIAMOND,A4} - 0.665 {}^0G_{Ni}^{FCC,A1} = -38227.151 + 4.849T$$

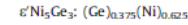


$${}^0G_{GeNiVa}^{\delta Ni_5Ge_3} - 0.5 {}^0G_{Ge}^{DIAMOND,A4} - 0.5 {}^0G_{Ni}^{FCC,A1} = -54286.304 - 5.624T$$

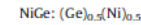
$${}^0G_{GeNiNi}^{\delta Ni_5Ge_3} - 0.5 {}^0G_{Ge}^{DIAMOND,A4} - 2 {}^0G_{Ni}^{FCC,A1} = -110540 + 11.717T$$

$${}^0G_{GeNiVa}^{\delta Ni_5Ge_3} = -2655.913 - 2.932T$$

$${}^1G_{GeNiVa}^{\delta Ni_5Ge_3} = -17558.144$$



$${}^0G_{Ni_5Ge_3}^{\epsilon' Ni_5Ge_3} - 0.375 {}^0G_{Ge}^{DIAMOND,A4} - 0.625 {}^0G_{Ni}^{FCC,A1} = -37350.646 + 3.328T$$



$${}^0G_{NiGe}^{\delta NiGe} - 0.5 {}^0G_{Ge}^{DIAMOND,A4} - 0.5 {}^0G_{Ni}^{FCC,A1} = -30992.547 + 0.967T - 0.17 \ln T + 6.015E - 05T^2 - 9.471E - 08T^3 + 2.393E - 22T^7 - 14960.491T^{-1}$$

^a Gibbs energies are expressed in J/mol. The lattice stabilities were given by Dinsdale [44].

^b The ordered phase βNi_3Ge with $L1_2$ structure and disordered phase FCC_A1 are modeled as the same phase.

^c The sublattice model of the Ni_5Ge_3 phase is given as integral stoichiometry since 1/3 cannot be exactly expressed as a decimal.

Thermodynamic modeling

Example: assessment of the binary Ge-Ni system

- Work flow for thermodynamic assessment
 - Comparison simulation vs experimental data

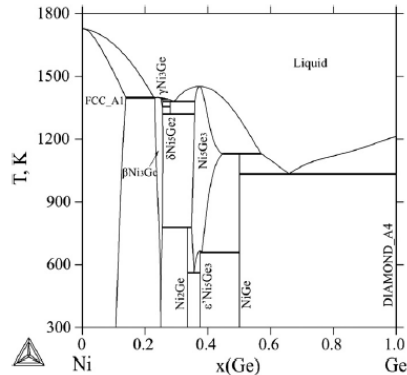


Fig. 4. Calculated phase diagram of the Ge-Ni binary system from the present work.

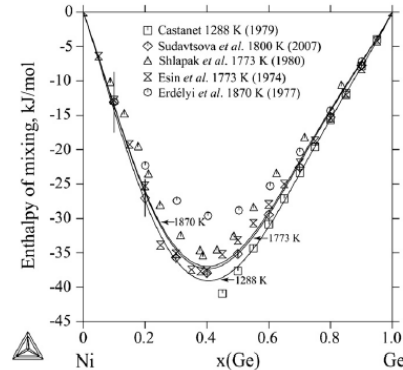


Fig. 6. Calculated enthalpy of mixing of liquid alloys compared with experimental data. (Ref. states: liquid Ge and liquid Ni).

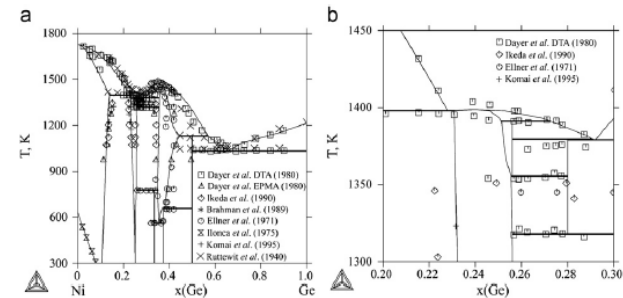


Fig. 5. (a) Calculated phase diagram of the Ge-Ni binary system compared with experimental data. (b) Ni-rich part of the phase diagram compared with experimental data.

Thermodynamic modeling

Example: assessment of the binary Ge-Ni system

- Work flow for thermodynamic assessment
 - Comparison simulation vs experimental data

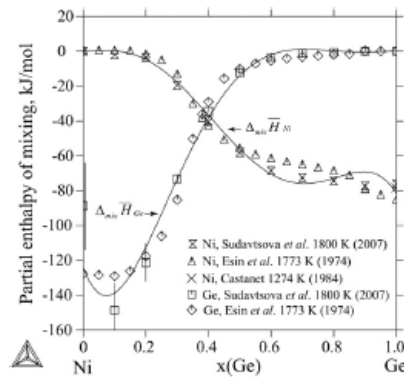


Fig. 7. Calculated partial enthalpy of Ni and Ge in liquid alloys at 1800 K compared with experimental data. (Ref. states: liquid Ge and liquid Ni).

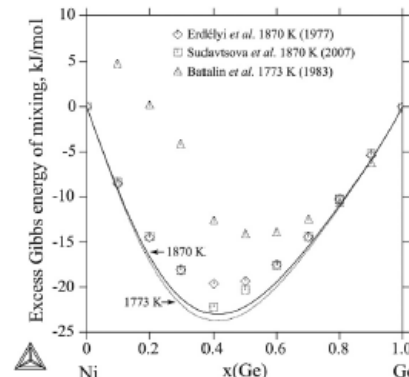


Fig. 9. Calculated excess Gibbs energy of mixing of liquid alloys compared with experimental data. (Ref. states: liquid Ge and liquid Ni).

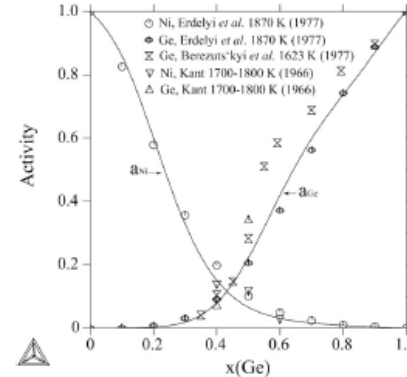


Fig. 8. Calculated activity of Ge and Ni in the Ge-Ni melt at 1870 K compared with experimental data. (Ref. states: liquid Ge and liquid Ni).

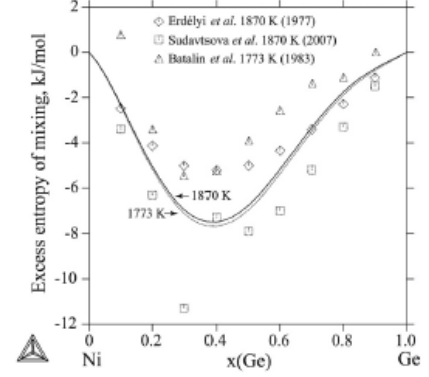


Fig. 10. Calculated excess entropy of mixing of liquid alloys compared with experimental data. (Ref. states: liquid Ge and liquid Ni).

Thermodynamic databases

Steel and Fe-alloys

TCFE11 Elements, Systems, Phases and Properties

Included Elements

There are 29 elements included in the most recent version of the database.

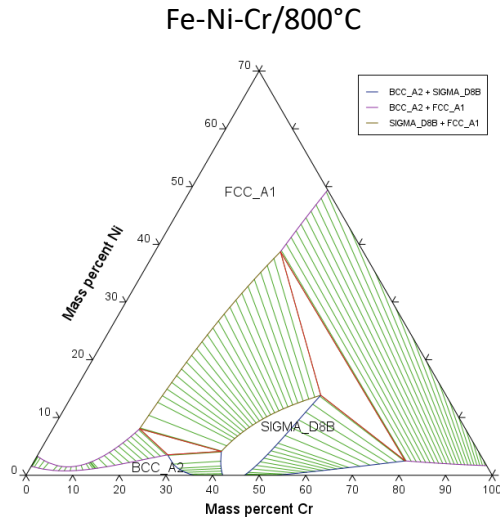
Ar	Al	B	C	Ca	Ce	Co	Cr	Cu	Fe
H	Mg	Mn	Mo	N	Nb	Ni	O	P	Ru
S	Si	Ta	Ti	V	W	Y	Zn	Zr	

The most recent version of the database contains the following:

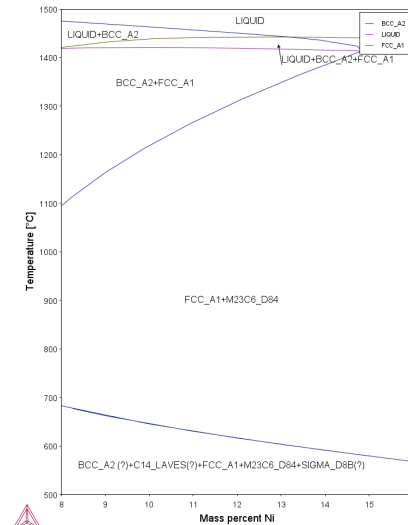
- 345 assessed binary systems
- 290 assessed ternary systems
- 79 assessed quaternary systems
- Several assessed quinary systems

Non-equilibrium transformations

- Thermodynamic simulation software tools allow calculating phase equilibria in multi-component systems as well as equilibrium thermodynamic properties, i.e. complete diffusion in the liquid and solid states are assumed



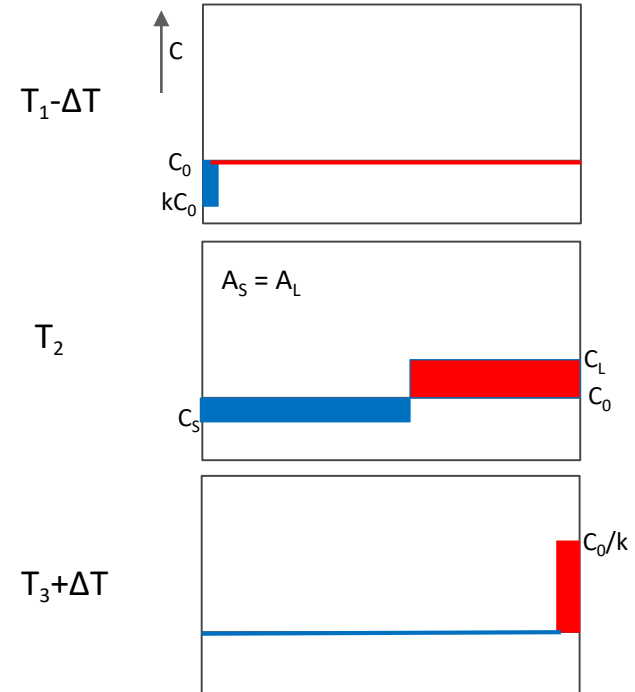
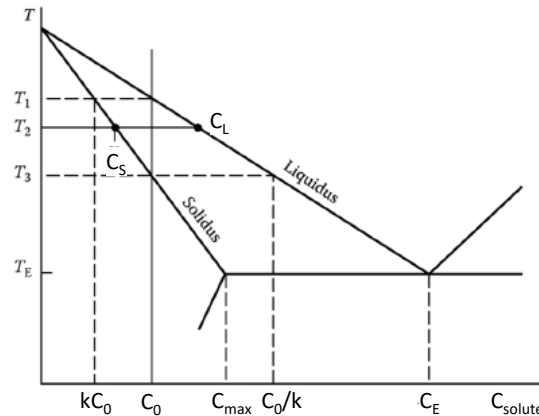
X5CrNiMo18-12-3, influence of Ni



- Thermodynamic simulation software tools allow calculating phase equilibria in multi-component systems as well as equilibrium thermodynamic properties, i.e. complete diffusion in the liquid and solid states are assumed
- Metals processing includes many non-equilibrium steps
 - casting & solidification
 - Solutionizing/homogenization heat treatments
 - Quenching and precipitation hardening
 - Joining (soldering/brazing, diffusion bonding)
- During service of high-performance alloys, diffusion-controlled processes can take place
 - Grain coarsening
 - Precipitate growth
 - Formation of intermetallic compounds at interfaces
- How can these non-equilibrium and time-dependent processes be modeled and simulated?

The equilibrium solidification model

- Under equilibrium conditions, the solidification path in an alloy is given by the lever rule
- The equilibrium solute concentrations are
 - $kC_0 \leq C_S \leq C_0$
 - $C_0 \leq C_L \leq C_0/k < X_E$



The Scheil-Gulliver solidification model

- In reality, solidification typically occurs under non-equilibrium conditions, i.e. the cooling rate is too high to allow time for complete redistribution of alloying elements according to equilibrium
- A qualitative description of the solute redistribution during solidification processes is possible with the so-called “Scheil–Gulliver” model, which was first formulated in 1913 by Gulliver.
- The basic assumptions of the model are:
 - Diffusion of all elements in the liquid phase is infinitely fast
 - Diffusion of all elements in the solid phases is zero
 - The liquid/solid interface is in thermodynamic equilibrium

The Scheil-Gulliver solidification model

- The corresponding differential equation for the mass balance between solid and liquid was presented by Scheil in 1942:

$$(C_L - C_S)df_s = (f_L)dC_L$$

where C_S is the local composition of the solid, C_L of the liquid, f_s is the fraction of solid and f_L is the fraction of liquid

- With the partitioning coefficient $k = C_S/C_L$ and considering that mass is conserved ($f_L + f_s = 1$), the mass balance may be rewritten as

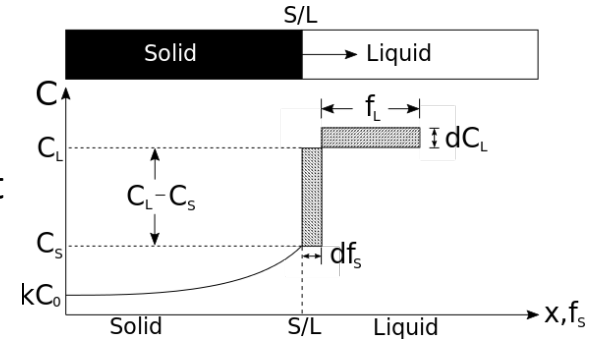
$$C_L(1 - k)df_s = (1 - f_s)dC_L$$

- Using the boundary condition $C_L = C_0$ at $f_s = 0$

$$\int_0^{f_s} \frac{df_s}{1 - f_s} = \frac{1}{1 - k} \int_{C_0}^{C_L} \frac{dC_L}{C_L}$$

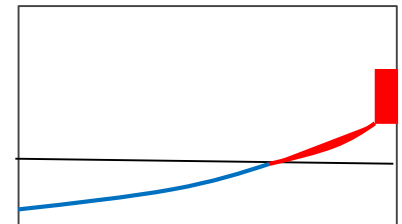
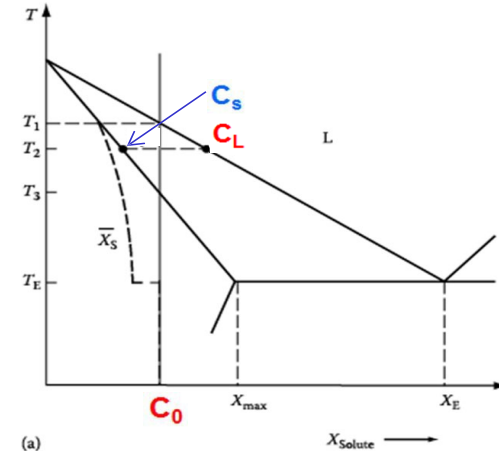
- Integrating results in the Scheil-Gulliver equation for composition of the liquid/solid during solidification gives

$$C_L = C_0(f_L)^{k-1}; C_S = kC_0(1 - f_s)^{k-1}$$

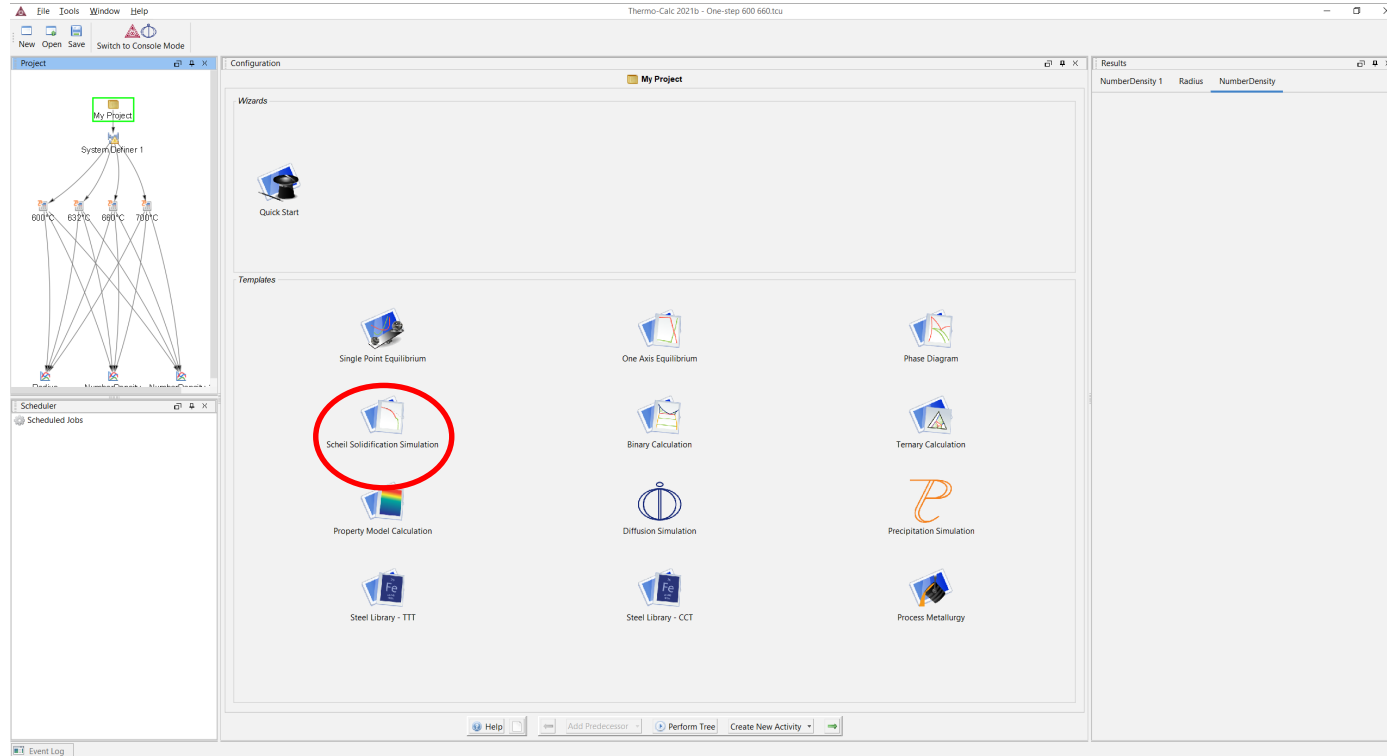


The Scheil-Gulliver solidification model

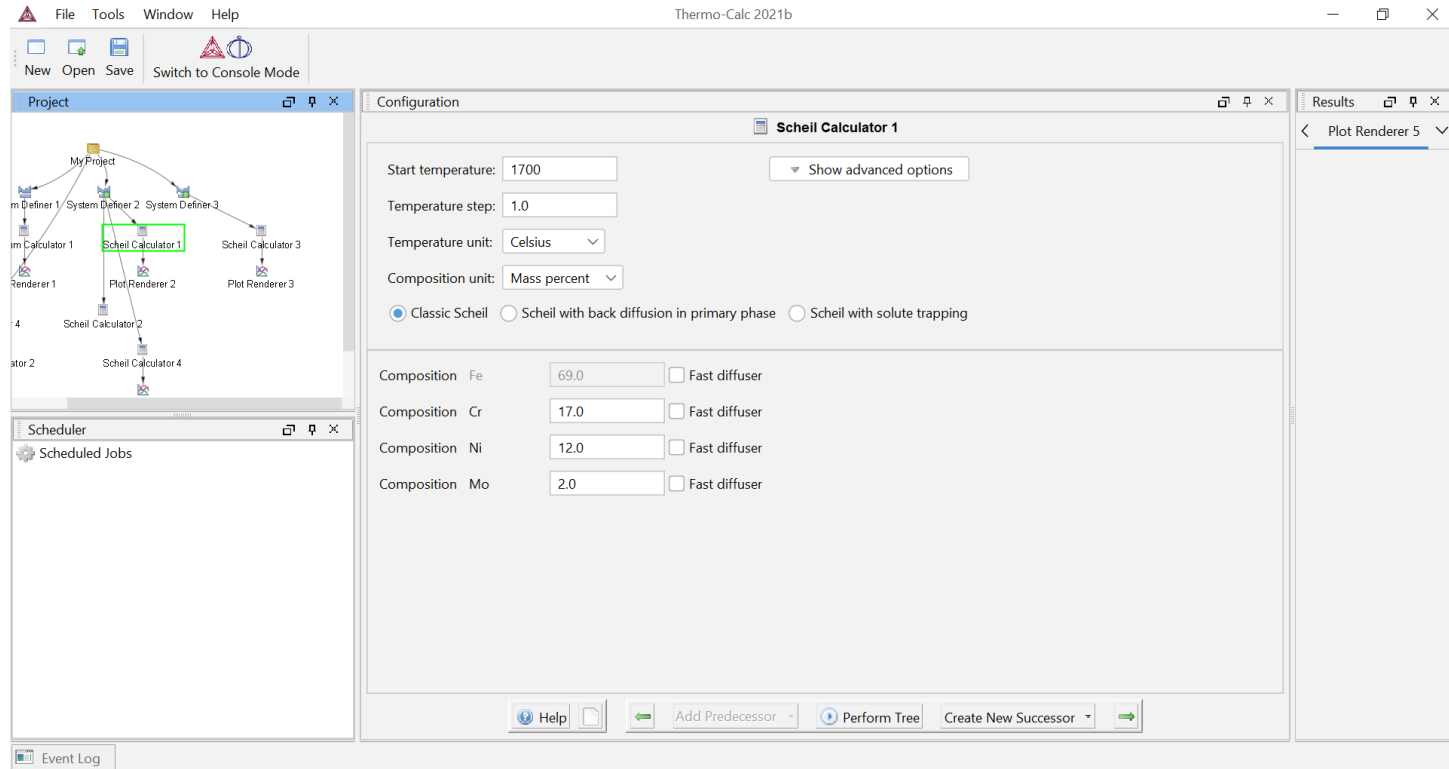
- Because of complete diffusion in the liquid, the liquid composition is always uniform, although this uniform composition changes as solidification continues
- The composition of the solid varies continuously as the solidification front advances and the solid retains this compositional variation after the front has passed
- The composition of the solid deviates from the equilibrium composition
- Solidification continues down to the eutectic temperature; when T_E is reached, the remaining liquid will have the eutectic composition, and it solidifies to form a eutectic solid.



The Scheil-Gulliver module

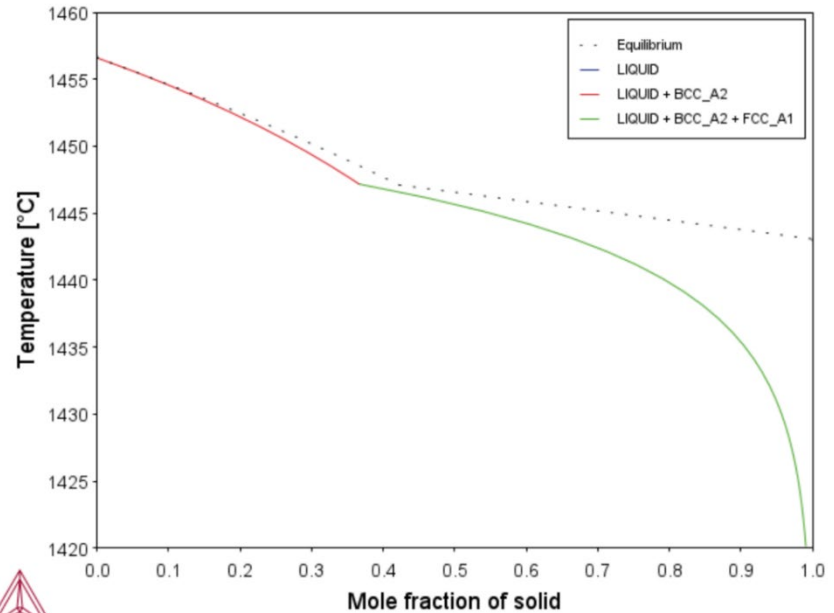
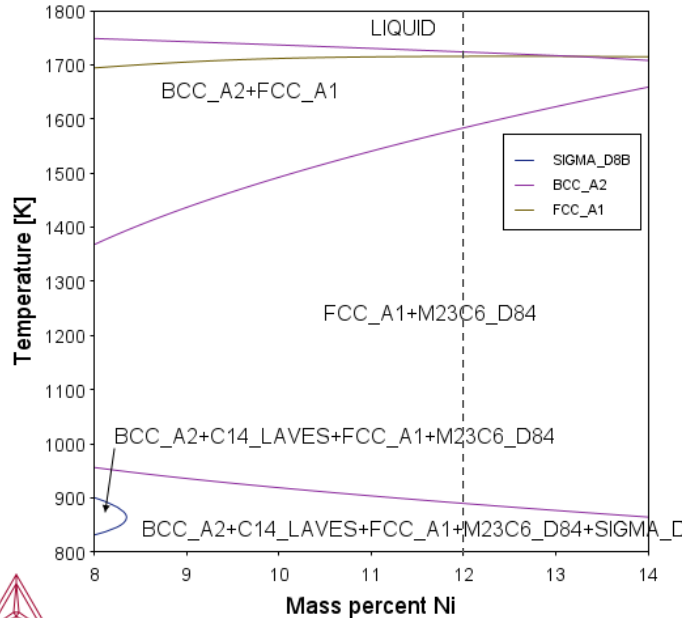


The Scheil-Gulliver module



The Scheil-Gulliver solidification model

Solidification of 1.4404/X3CrNiMo 17-12-2



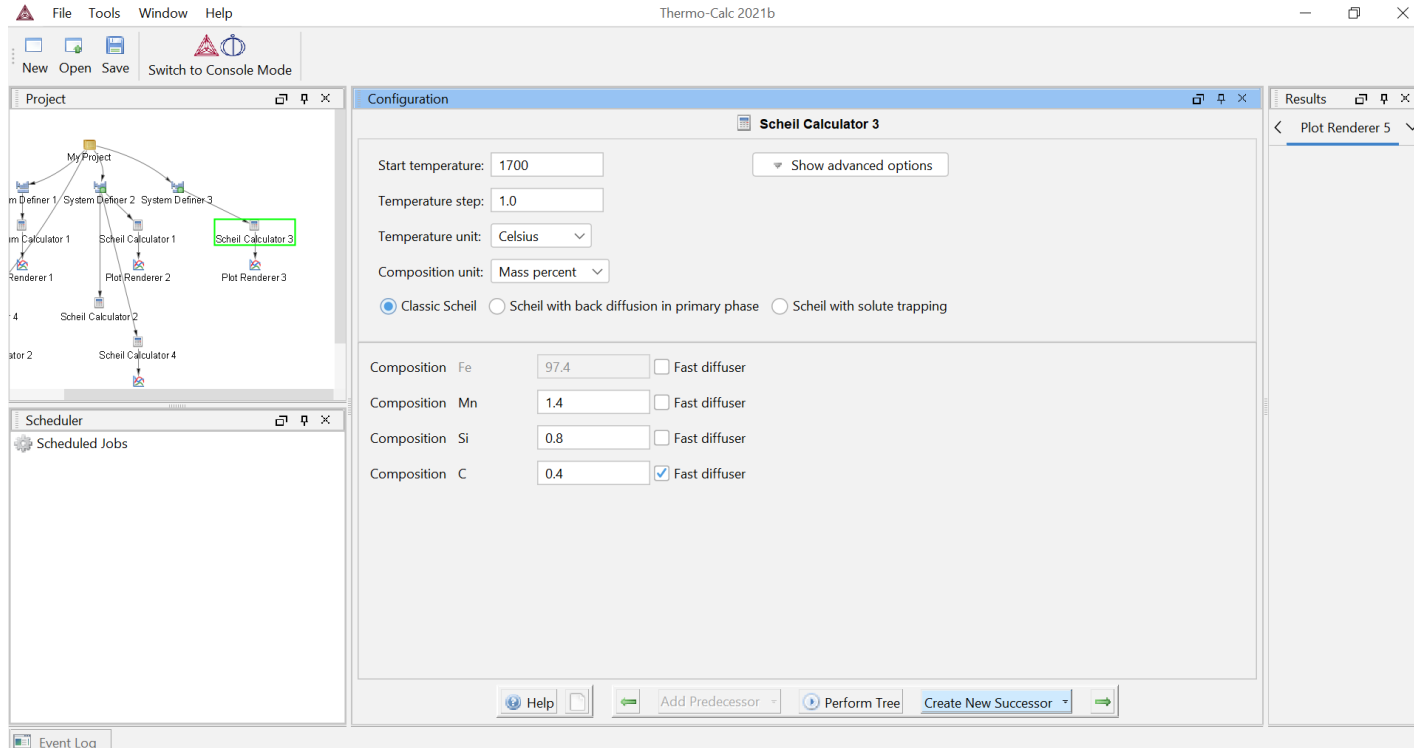
Modified Scheil-Gulliver models

S-G with fast diffusers

- Interstitial elements in steel like C, N or O have high diffusion rates
- The assumption of no diffusion in the solid phase during solidification is thus not correct at most industrial solidification rates.
- Scheil with fast diffusers is a variant of the classic Scheil simulation that was developed mainly for steel applications
- The assumptions are
 - Diffusion of all elements in the liquid phase is infinitely fast
 - Diffusion of all elements in the solid phases except the ones defined as “fast diffusers” is zero
 - Diffusion of the elements defined as “fast diffusers” is infinitely fast in the solid phase
 - The liquid/solid interface is in thermodynamic equilibrium

Modified Scheil-Gulliver models

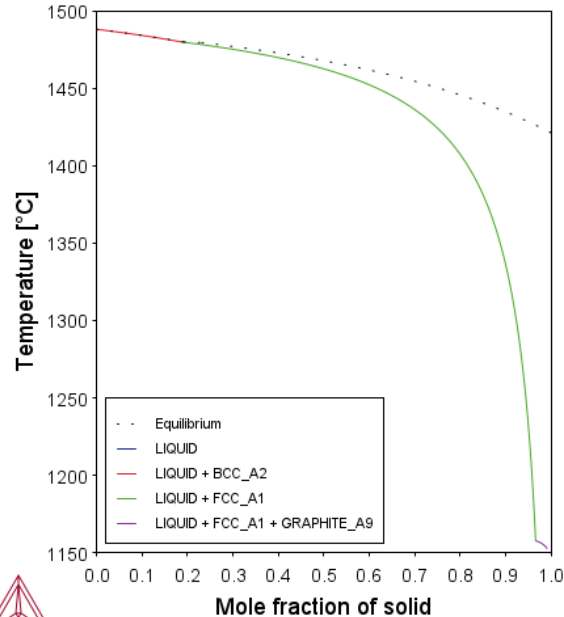
S-G with fast diffusers



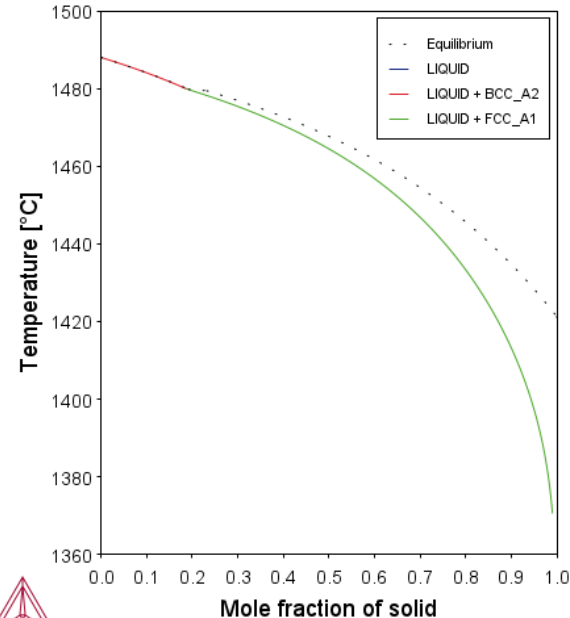
Modified Scheil-Gulliver models

Solidification of Fe-1.4Mn-0.8Si-0.4C

Classical Scheil



Modified Scheil
with C as fast diffuser



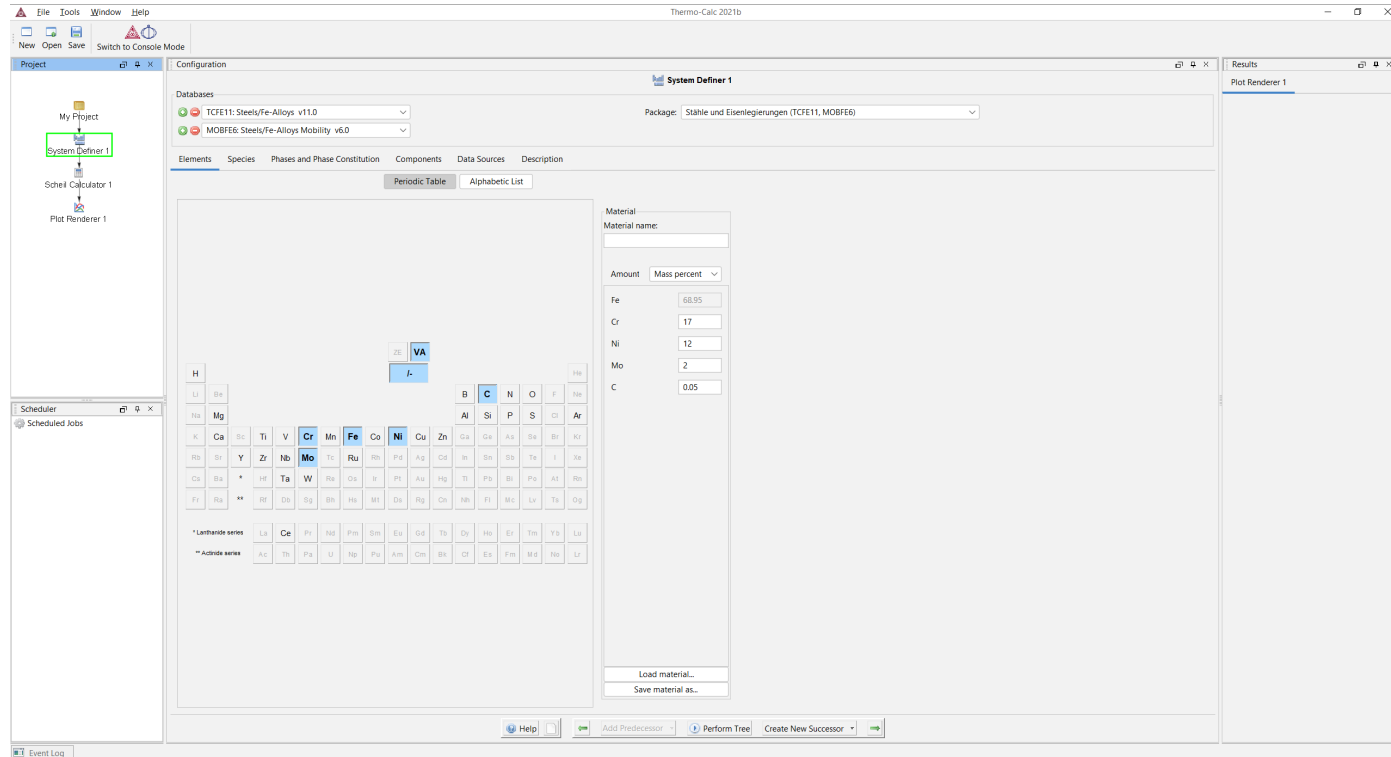
Modified Scheil-Gulliver models

S-G with back-diffusion in primary phase

- This model uses diffusion data from a mobility database to calculate back diffusion in the primary phase
- As with the Scheil-Gulliver model, this model assumes infinitely fast diffusion in the liquid material, but it allows for limited diffusion in the primary phase of the solid material.
- The assumptions are
 - Diffusion of all elements in the liquid phase is infinitely fast
 - Diffusion of all elements in the primary solid phase is quantitatively calculated using mobility data, a cooling rate, and a domain size (typically this will be the secondary arm spacing)
 - The liquid/solid interface is in thermodynamic equilibrium
- Using diffusion data from a mobility database, it quantitatively takes into account the real back diffusion of all elements in the primary solid phase (typically the FCC or BCC phase)

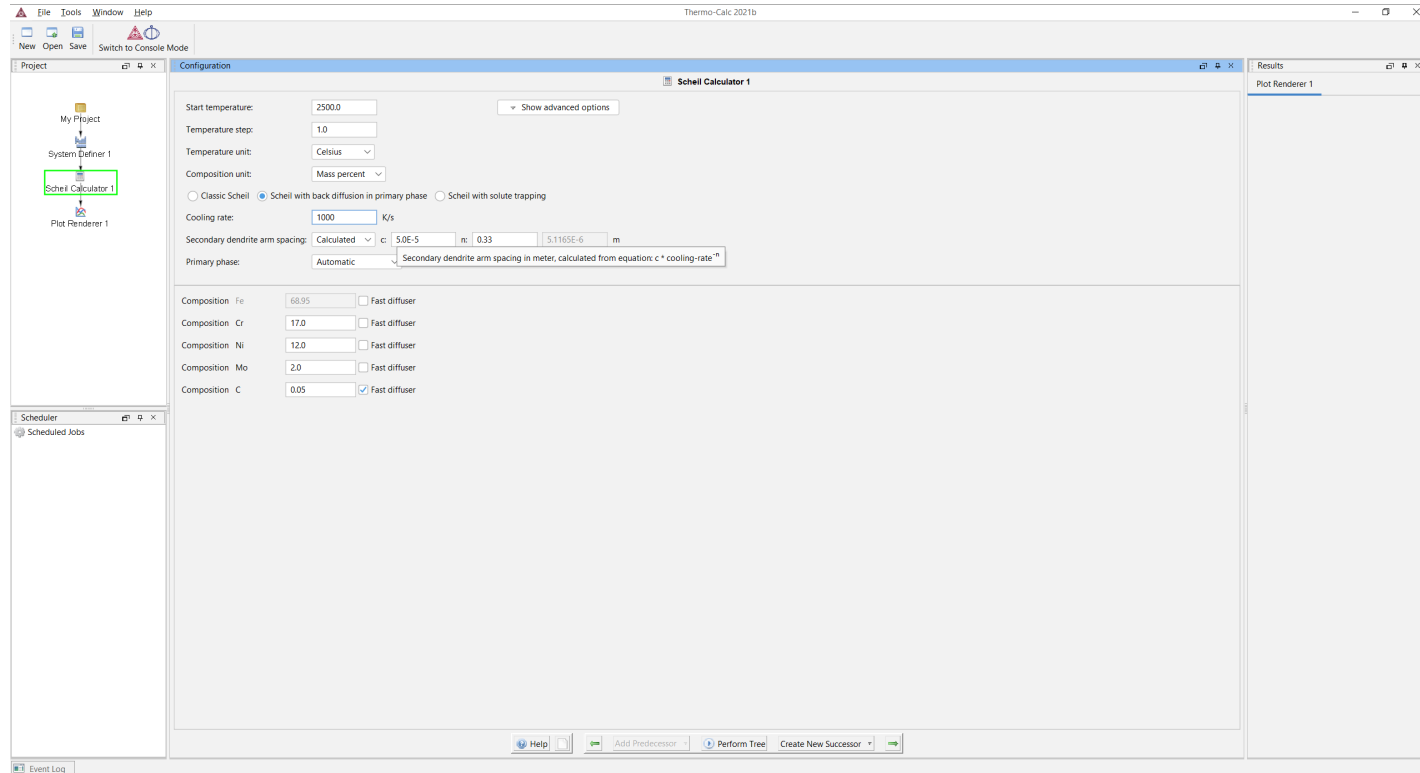
Modified Scheil-Gulliver models

S-G with back-diffusion in primary phase



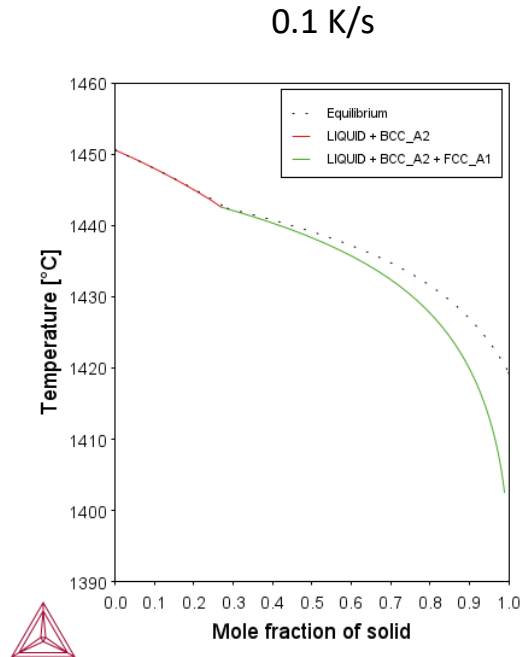
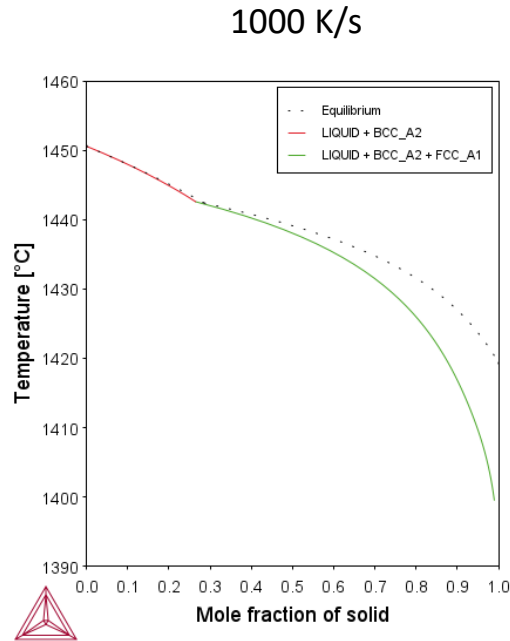
Modified Scheil-Gulliver models

S-G with back-diffusion in primary phase



Modified Scheil-Gulliver models

Solidification of 1.4404/X3CrNiMo 17-12-2



- Fick's 1st and 2nd law describe the flux of a species k in a concentration gradient (1) and the solute balance as a function of time (2), respectively

$$J_k = -D_k \frac{\partial c_k}{\partial x} \quad (1) \text{ and } \frac{\partial c_k}{\partial t} = D_k \frac{\partial^2 c_{Bk}}{\partial x^2} \quad (2)$$

- If an alloy contains 2 or more species, the diffusivities D_k not only depend on the concentration but also on the concentration gradient
- A multicomponent extension to Fick's 1st law was first expressed by Onsager in 1931, postulating that each flux is linearly related to every thermodynamic force

$$J_k = - \sum_{i=1}^n L'_{ki} \frac{\partial \mu_i}{\partial x}$$

where the μ_i terms are the chemical potentials for the various species, which are unique functions of the composition ($\mu_i = f(c_1, c_2, c_3, \dots, c_n)$) and L'_{ki} is the proportionality factor, which depends on the mobility of the individual species

- The flux J_k is defined such as $\sum_{k=1}^n V_k J_k = 0$ (V_k : partial molar volume of k)

- Generally, it is convenient to express the fluxes as functions of gradients in concentration rather than gradients in chemical potential

$$J_k = - \sum_{i=1}^n L'_{ki} \sum_{j=1}^n \frac{\partial \mu_i}{\partial c_j} \frac{\partial c_j}{\partial x}$$

- With $D_{kj} = - \sum_{i=1}^n L'_{ki} \frac{\partial \mu_i}{\partial c_j}$ (unreduced diffusivities)

$$J_k = - \sum_{j=1}^n D_{kj} \frac{\partial c_j}{\partial x}$$

- The $\frac{\partial \mu_i}{\partial c_j}$ values are the so-called thermodynamic factors
→ the diffusivities consist of a thermodynamic and a kinetic part

- Atomic mobilities are purely kinetic, element specific properties, defined as the velocity of a species per unit force (unit: $\frac{m^2/s}{J/mol}$)

- The mobility coefficient for an element B in metallic matrix can be written as

$$M_B = \frac{M_B^0}{RT} \exp\left(\frac{-Q_B}{RT}\right)$$

where M_B^0 is the compositional-dependent frequency factor and Q_B the activation enthalpy

- The composition-dependency can be represented as a linear combination of the values at each endpoint of the composition space and a Redlich-Kister term

$$\Phi_B = \sum_i x_i \Phi_B^i + \sum_i \sum_{j>i} x_i x_j \left[\sum_{r=0}^m {}^r A_B^{i,j} (x_i - x_j)^r \right]$$

Φ_B represents $\ln M_B^0$ or Q_B

Φ_B^i is the value of Φ_B for pure i and thus represents one of the endpoints in the composition space

${}^r A_B^{i,j}$ are binary interaction parameters, the commas separating different species interacting with each other

x_i and x_j are mole fractions for elements i and j, respectively

- As in the case of thermodynamic data, the model parameters $\Phi_B^i, {}^r A_B^{i,j}, \dots$ are determined by an optimization procedure (considering experimental data)

- The multi-component diffusion theory has been implemented in DICTRA (now TC diffusion module)
- The mobilities are stored in the mobility databases
- The reason to store individual mobilities rather than interdiffusion coefficients is that in an n-component system, there are n mobilities and $(n-1)^2$ interdiffusion coefficients
- The mobilities are related to the interdiffusion coefficients according to

$$L'_{ki} = \sum_{i=1}^n (\delta_{ik} - c_k V_i) c_i y_{va} M_i$$

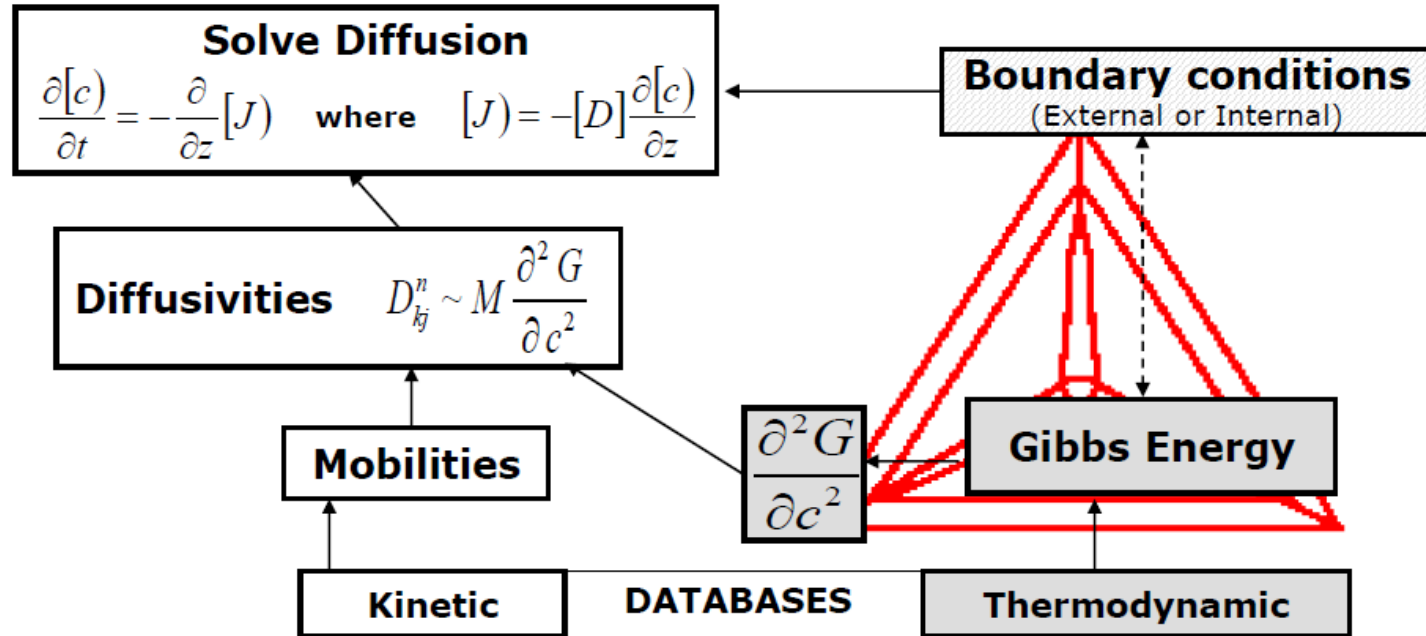
δ_{ik} is the Kronecker delta, i.e. $\delta_{ik} = 1$ when $j=k$ and $\delta_{ik} = 0$ otherwise

c_k and c_i are the amounts of k and i per unit volume

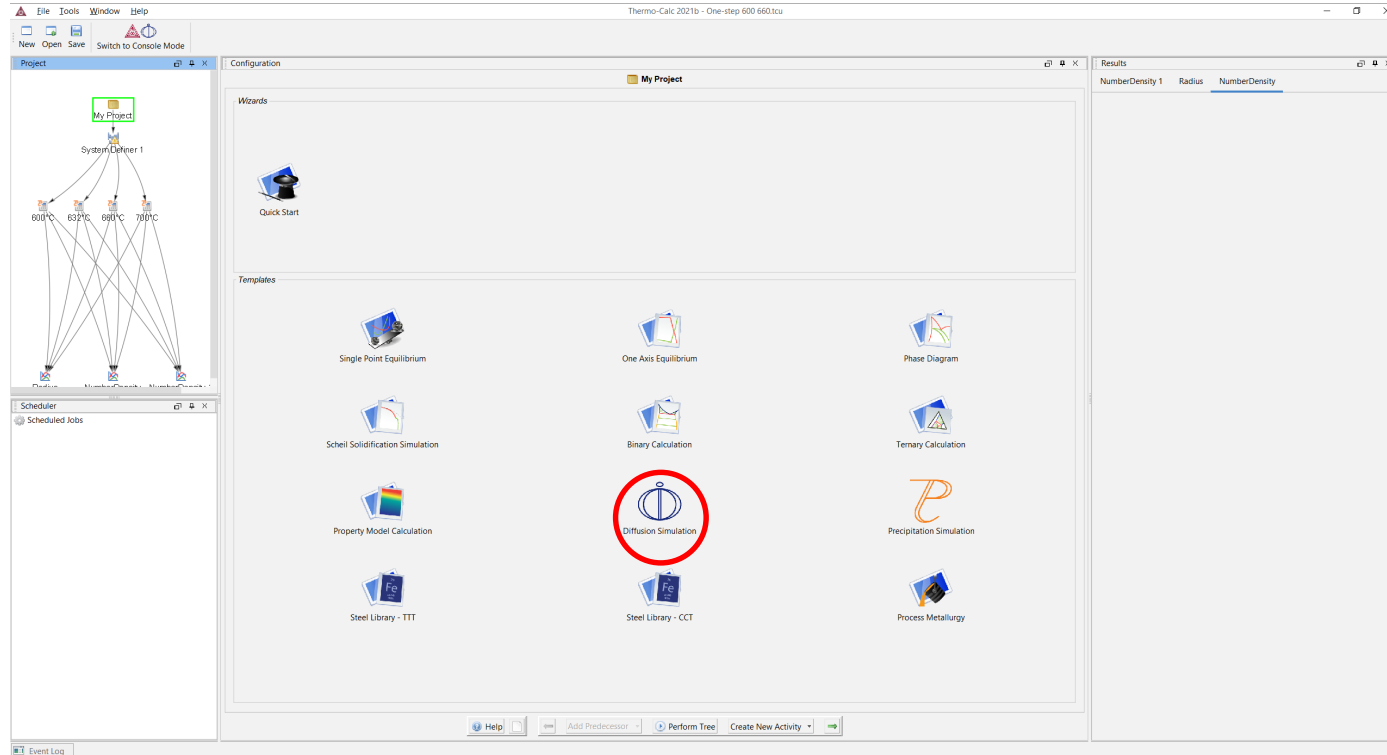
V_i is the partial molar volume of element i

y_{va} is the fraction of vacant lattice sites on the sublattice where i is dissolved

M_i is the mobility of i when i is interstitial and the mobility divided by y_{va} when i is substitutionally dissolved



The diffusion simulation module



The diffusion simulation module

Example: Brazing of diamond (C) to steel

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Modeling and simulation of the TiC reaction layer growth during active brazing of diamond using DICTRA



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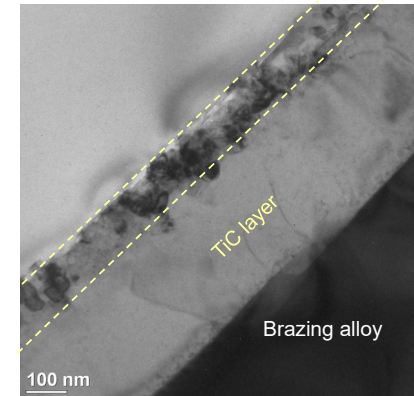
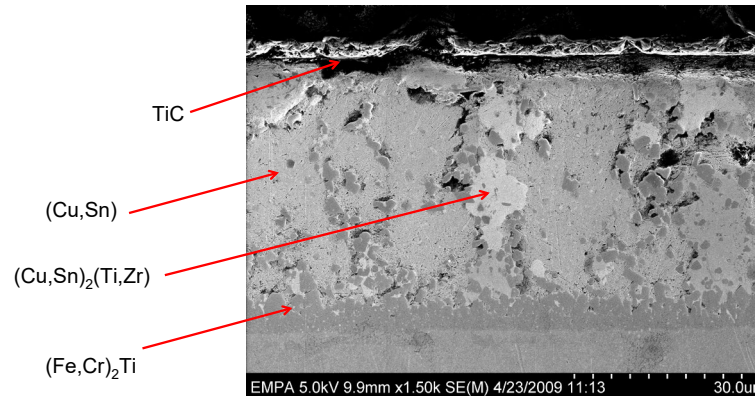
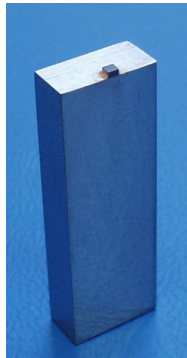
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The diffusion simulation module

Example: Brazing of diamond (C) to steel

- Vacuum brazing 930 °C /10 min
- Filler: Cu-14.4Sn-10.2Ti-1.5Zr (mass %)
- Base material: X2Cr Ni Mo 18 14 3 stainless steel (AISI 316)
- Thickness of brazing layer: ~50 μm



The diffusion simulation module

Example: Brazing of diamond (C) to steel

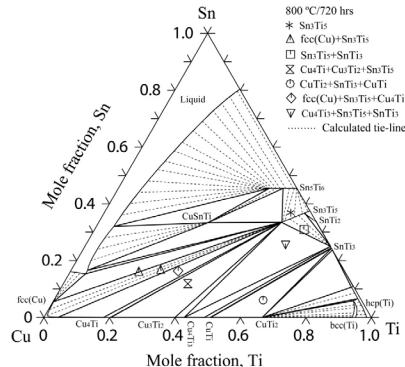
Thermodynamics

$$G_{mix} = G^0 + G_{mix}^{ideal} + G_{mix}^{excess}$$

$$G_{mix} = \sum_i x_i G_i^0 + RT \sum_i x_i \ln x_i + \sum_i \sum_{j>i} x_i x_j \Omega_{ij}$$

$$\Omega_{i,j}^{\phi} = {}^0\Omega_{i,j}^{\phi} + {}^1\Omega_{i,j}^{\phi}(x_i - x_j) + ({}^1\Omega_{i,j}^{\phi})^2(x_i - x_j)^2 \dots$$

$$k\Omega_{i,j}^{\phi} = A + BT + CT \ln T$$



Diffusion equation

Diffusivity

Thermodynamic factor

Mobility

Kinetics/Diffusion

$$J_k = -\sum_{i=1}^n D_{kj}^n \nabla \cdot c_i \quad \frac{\partial c_k}{\partial t} = -\nabla \cdot J_k$$

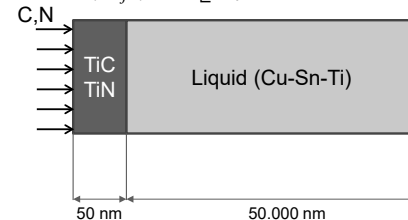
$$D_{kj}^n = \sum_{i=1}^n (\delta_{ik} - x_k) x_i M_i \left(\frac{\partial \mu_i}{\partial x_j} - \frac{\partial \mu_i}{\partial x_n} \right)$$

$$M_B = \exp\left(\frac{-Q_B + RT \ln M_B^0}{RT}\right) \frac{1}{RT} \text{ mg } \Omega$$

$$= \exp\left(\frac{\Delta G^*}{RT}\right) \frac{1}{RT} \text{ mg } \Omega$$

$$\Delta G^* = \Phi_B$$

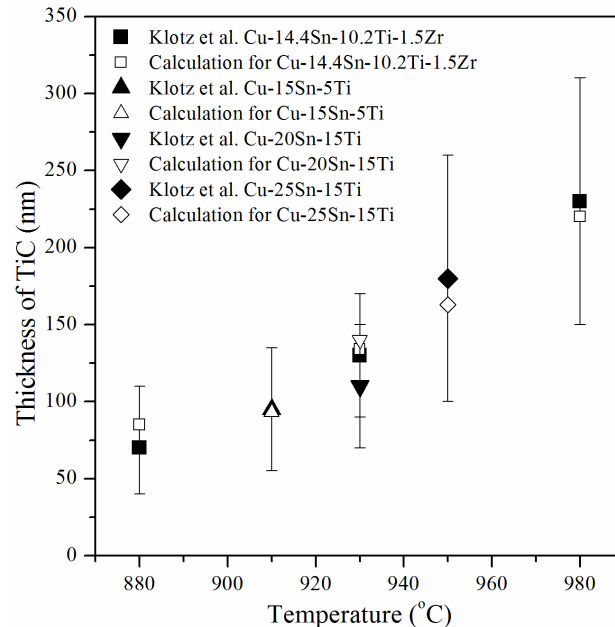
$$= \sum_i x_i \Phi_B^i + \sum_i \sum_{j>i} x_i x_j \left[\sum_{r=0}^m {}^r \Phi_B^{i,j} (x_i - x_j)^r \right]$$



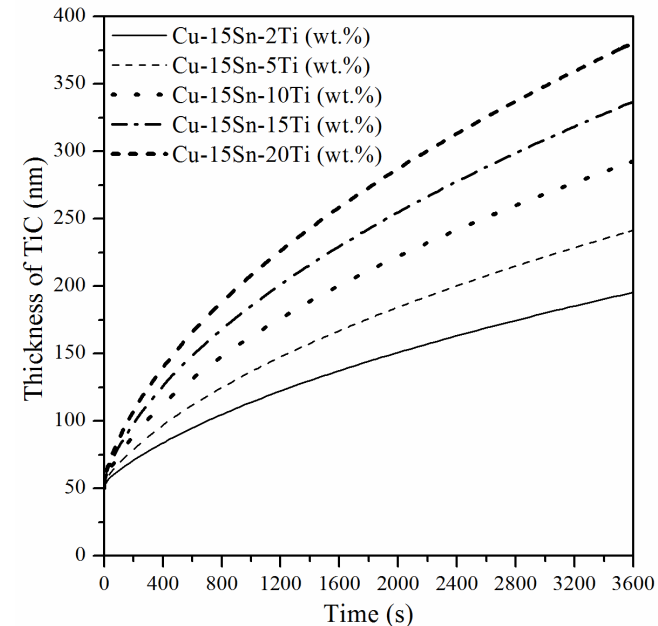
The diffusion simulation module

Example: Brazing of diamond (C) to steel

- Simulated influence of alloy composition + experimental results



- Simulated influence of amount of Ti on TiC growth in brazing diamond at 930°C

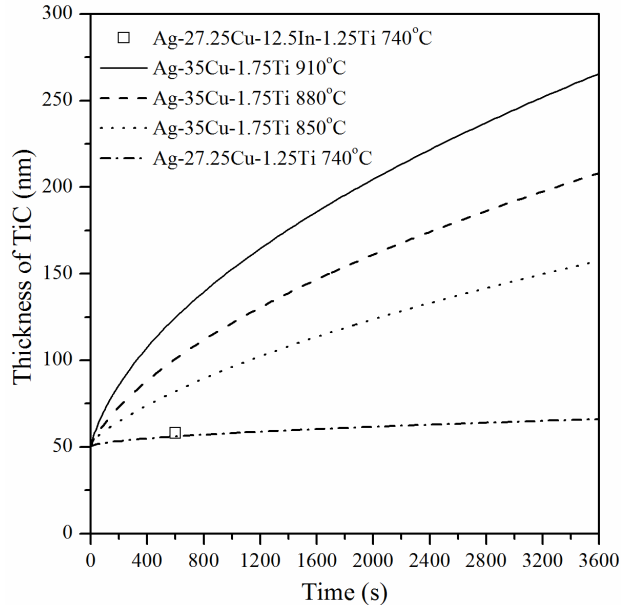


/W.J. Zhu, C. Leinenbach et al, Comp. Mater. Sci. 78 (2013) 74./

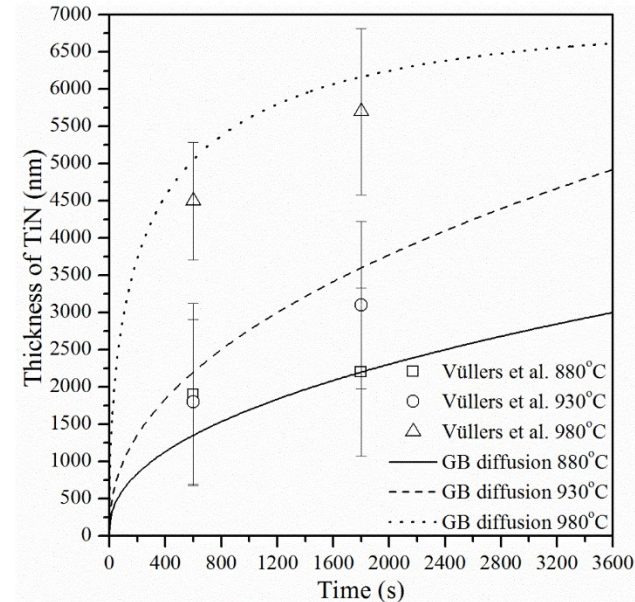
The diffusion simulation module

Example: Brazing of diamond (C) to steel

- Simulated influence of Ag-based filler alloy



- TiN reaction layer growth during brazing of AlN ceramics

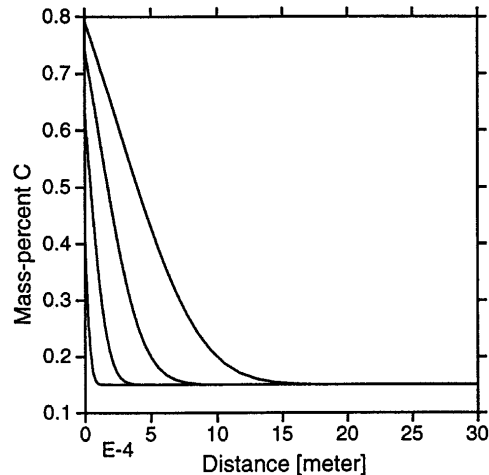


/W.J. Zhu, C. Leinenbach et al, Comp. Mater. Sci. 78 (2013) 74./

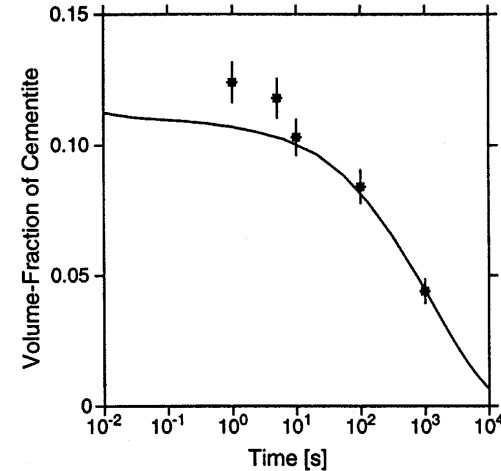
The diffusion simulation module

Other examples

Calculated C-concentration profiles after 100, 1000, 5'000 and 18'000 s at 900°C in a carburizing atmosphere



Calculated volume fraction of cementite as a function of time, during austenitization of an Fe-2.06 at.% Cr-3.91 at.% C alloy at 910°C



/A. Borgenstam et al., Journal of Phase Equilibria 21(3) (2000)

Simulation of precipitation reactions

- Phase transformations start with the formation of stable nuclei of a size larger than the critical radius grown at expenses of available elements in the matrix.
- The growth of the nucleus ends when the thermodynamic equilibrium is reached, i.e. chemical potentials of each element are equal anywhere in the system.
- Assuming spherical nuclei, the nucleation rate can be described as

$$J(t) = J_s \exp\left(-\frac{\tau}{t}\right)$$

J_s : steady state nucleation rate

τ : incubation time

$$J_s = Z\beta^* N_0 \exp\left(-\frac{\Delta G^*}{kT}\right)$$

Z : Zeldovich factor

β^* : rate at which atoms are attached to critical nucleus

N_0 : number of available nucleation sites per unit volume

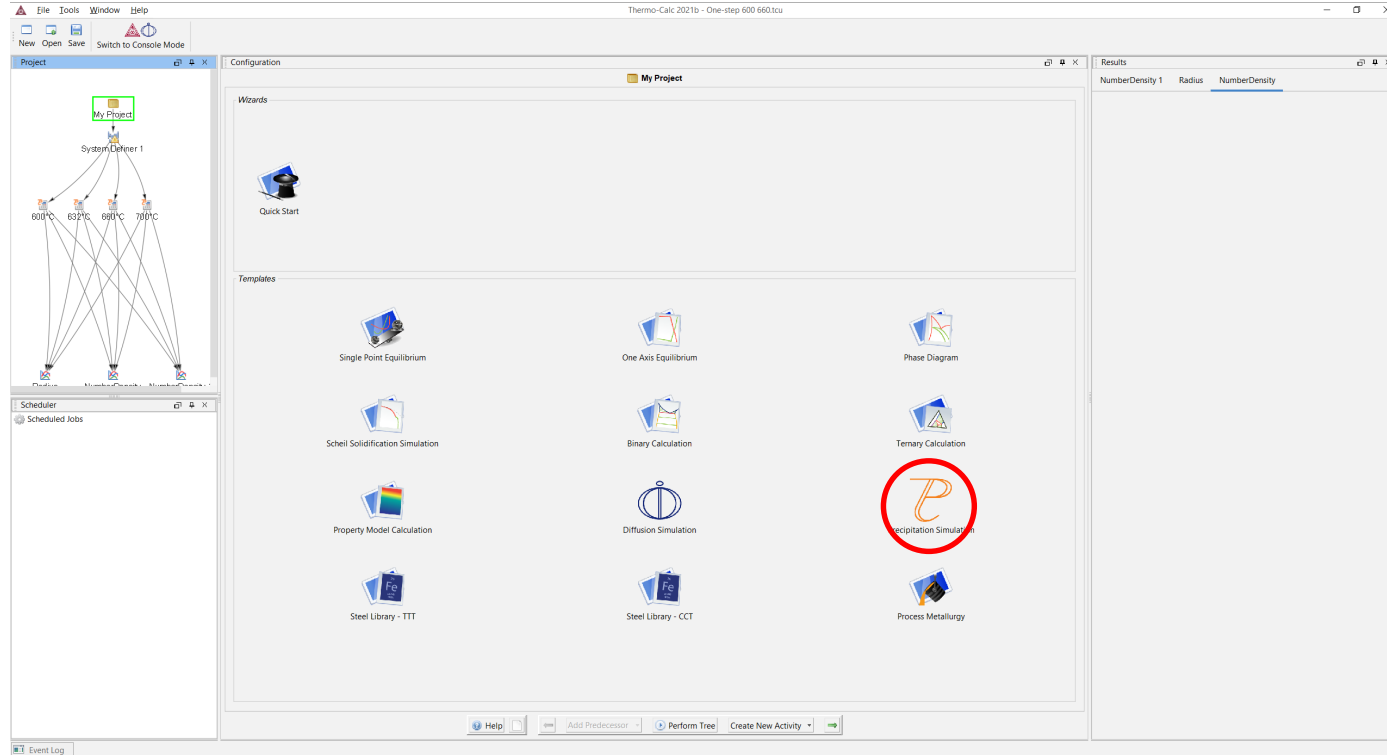
ΔG^* : nucleation barrier

$$\Delta G^* = \frac{16\pi\sigma^3}{3\Delta G_V^{\alpha \rightarrow P}}$$

σ : interfacial energy

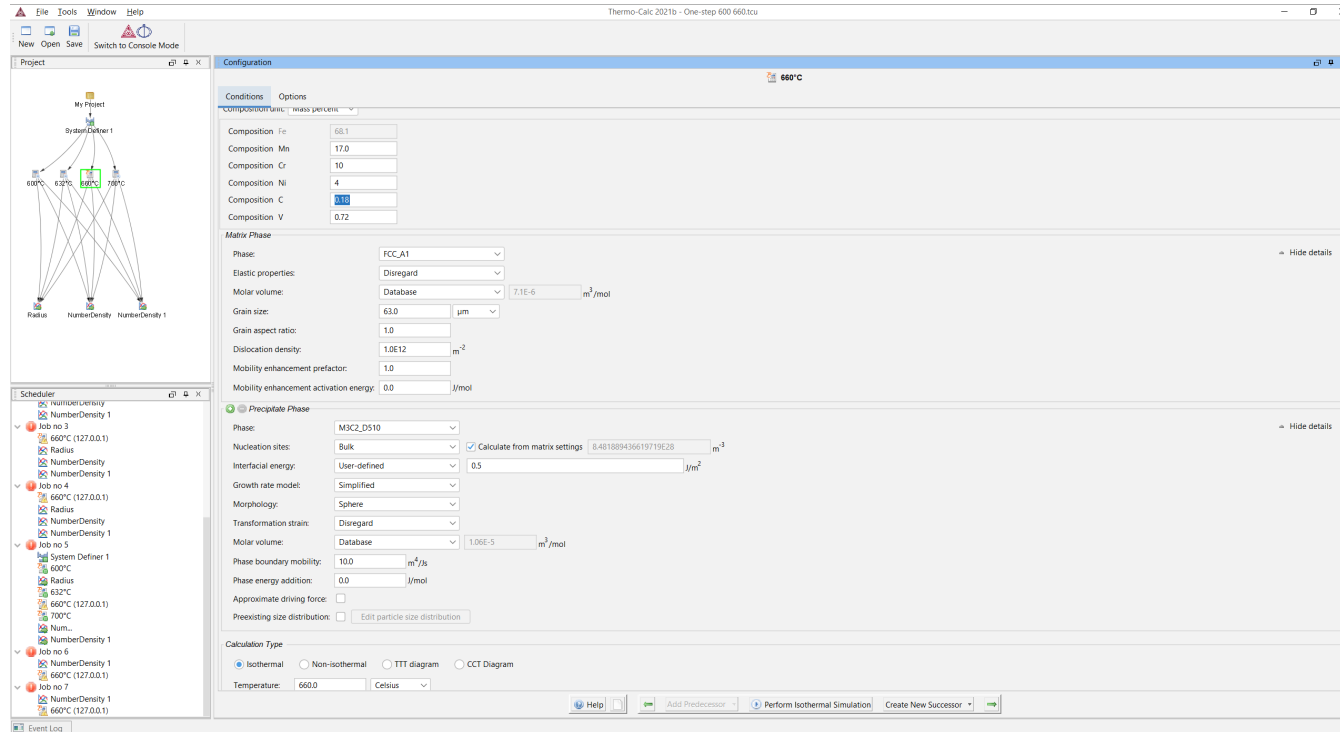
$\Delta G_V^{\alpha \rightarrow P}$: driving force for formation of precipitates P from α matrix

The precipitation module



The precipitation module

Example: Carbide precipitation in Fe shape memory alloy



- Thermodynamic and kinetic modeling
 - Understand the key concepts of thermodynamic modeling, including solution phases and the mathematical description of phase properties.
 - Differentiate between ideal, regular, and real solution models and their respective assumptions.
 - Redlich-Kister formalism for mathematical description of interactions
 - Understanding the concepts of advanced solution models: (sub)n-regular solution model; sub-lattice model; quasi-chemical model
 - Understanding the concepts of equilibrium and non-equilibrium solidification models:
 - Lever rule for equilibrium solidification
 - Scheil-Gulliver solidification for rapid cooling scenarios
 - Advanced variations of Scheil-Gulliver models (with fast diffusers, e.g., interstitial elements; with back-diffusion in primary solid phases).
- Practical applications
 - Use Thermo-Calc to assess thermodynamic and kinetic behaviors in multicomponent systems.
 - Casting and solidification, heat treatments (e.g., solutionizing, quenching, and precipitation hardening), joining processes (e.g., brazing, soldering, and diffusion bonding), grain growth